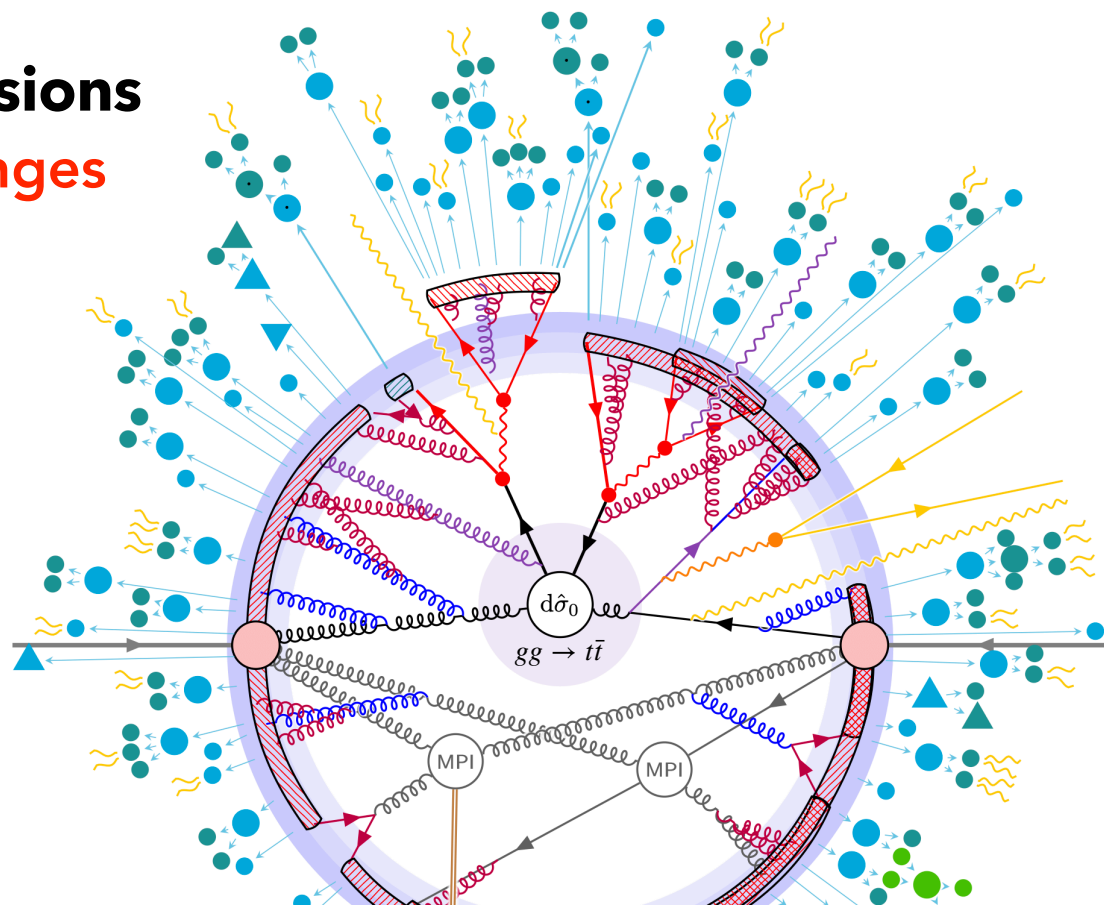


Anatomy of LHC collisions — and Future Challenges

Peter Z Skands

University of Oxford
& Monash University

IOP Meeting
Liverpool, April 2024



Australian Government
Australian Research Council



Anatomy of LHC collisions

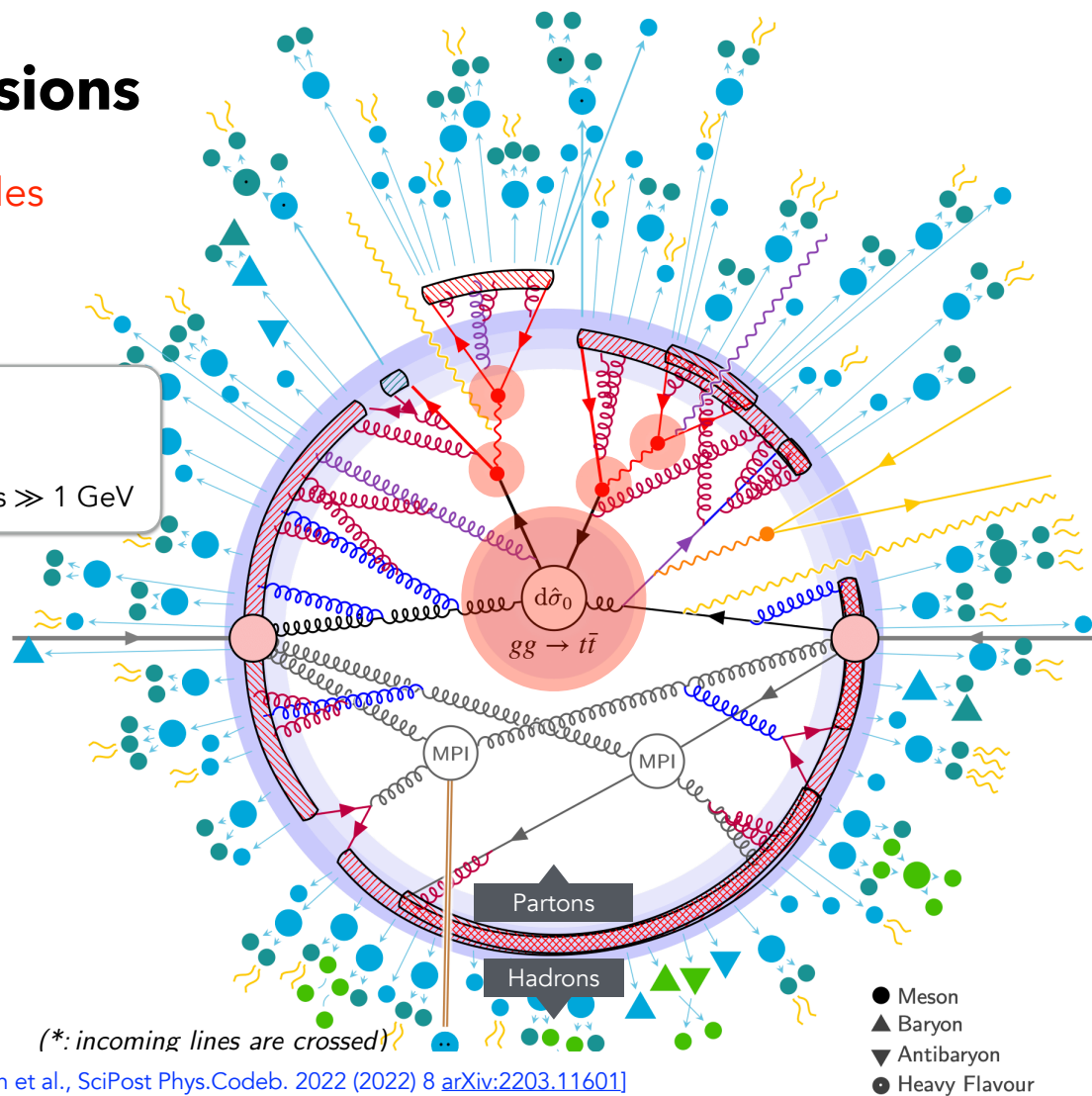
Physics Separation of scales

Maths Factorizations

Hard Processes

○ Hard Interaction
● Resonance Decays

"Hard" means large momentum transfers $\gg 1$ GeV



[Figure from Bierlich et al., SciPost Phys.Codeb. 2022 (2022) 8 arXiv:2203.11601]

Anatomy of LHC collisions

Physics Separation of scales

Maths Factorizations

Hard Processes

- Hard Interaction
- Resonance Decays

"Hard" means large momentum transfers $\gg 1$ GeV

Bremsstrahlung

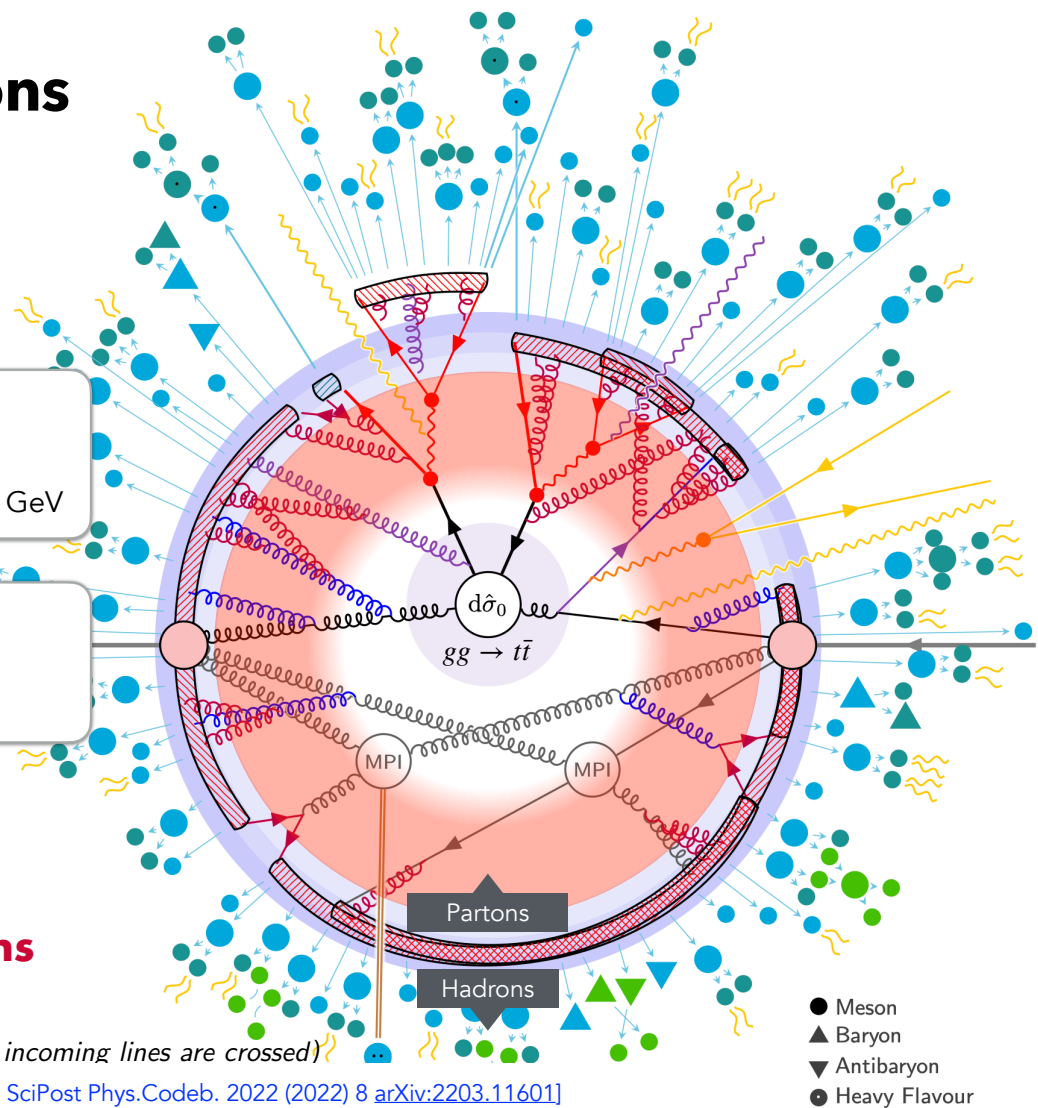
- Final-State Radiation
- Initial-State Radiation*

Down to momentum transfers ~ 1 GeV

For scales $\mu > 1$ GeV:

QCD Running Coupling $\alpha_s(\mu) < 1$

→ Perturbative **Approximations**



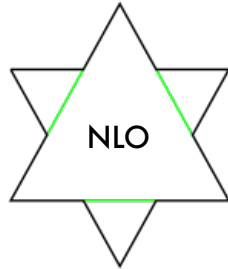
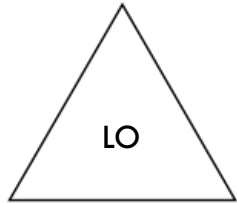
(*: incoming lines are crossed)

Perturbative Approaches

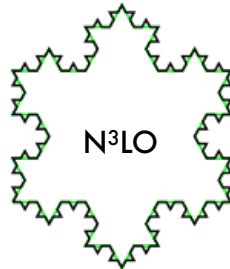
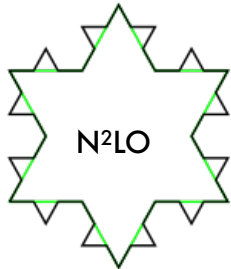
P.T. ~ Calculate the area of a shape ($d\sigma$) with higher and higher detail

Difference from exact area $\propto \alpha^{n+1}$

Fixed Order



Example: Koch Snowflake



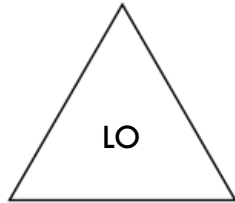
Note: (over)simplified analogy, mainly for IR structure. More at each order than shown here.

Perturbative Approaches

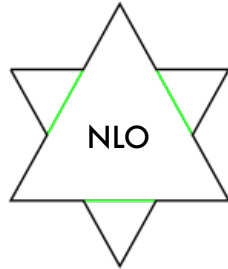
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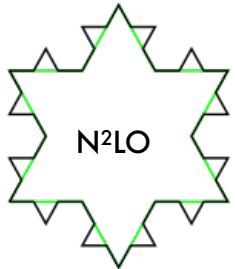


LO

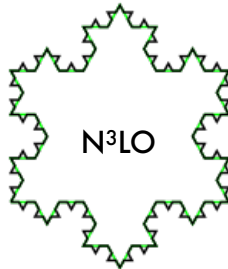


NLO

Example: Koch Snowflake

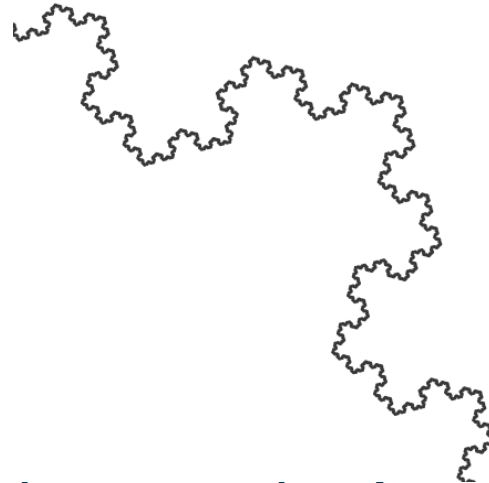


N²LO



N³LO

Resummation / Parton Showers



PROCTOR.CO

Massless gauge theories

Scale invariance → fractal substructure
(+ not hard to build in running coupling, masses)

Note: (over)simplified analogy, mainly for IR structure. More at each order than shown here.

Why go beyond **Fixed-Order** perturbation theory?

Schematic example:

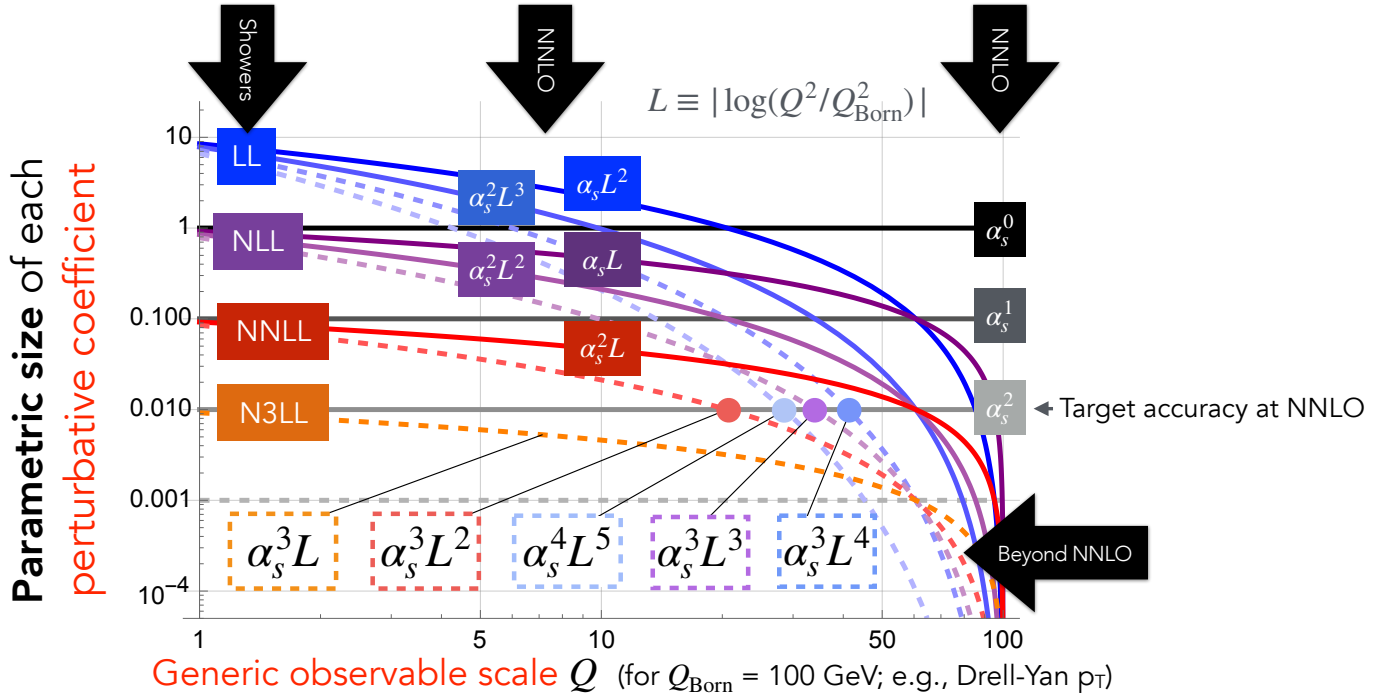
Calculation of the fraction of events passing a **radiation (jet) veto**:

$$\overbrace{\widehat{1}}^{\text{LO}} - \overbrace{\alpha_s(L^2 + L + F_1)}^{\text{NLO}} + \overbrace{\alpha_s^2(L^4 + L^3 + L^2 + L + F_2)}^{\text{NNLO}} + \dots$$

$$L \propto \ln(Q_{\text{veto}}^2 / Q_{\text{hard}}^2)$$

(Logs arise from integrals over propagators $\propto \frac{1}{q^2}$)

The Case for Embedding Fixed-Order Calculations within Showers



Bremsstrahlung Resummations (Showers) extend domain of validity of perturbative calculations

%-level precision @ LHC

⇒ NNLO + NNLL

= Our Target

Not quite there (yet) — but close ...

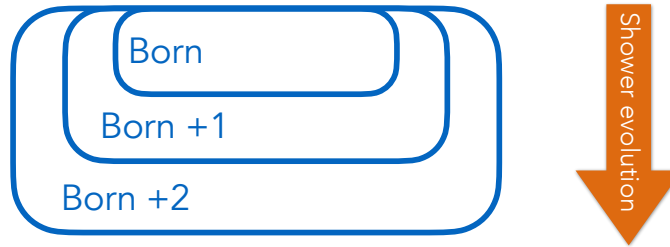
Towards True* NNLO Matching

*In the sense of the fixed-order and shower calculations matching each other point by point in each phase space

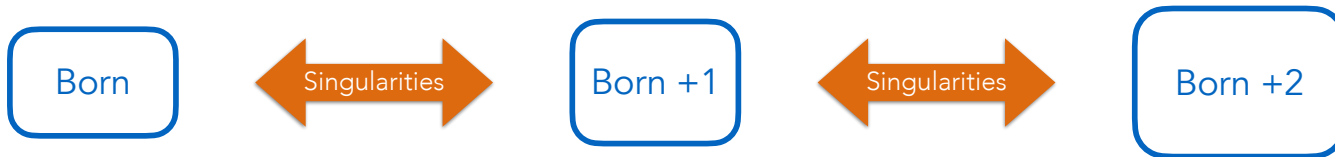
Idea: Use (nested) Shower Markov Chain as NNLO Phase-Space Generator

Harnesses the power of showers as efficient phase-space generators for QCD

Efficient: Pre-weighted with the (leading) QCD singular structures = soft/collinear poles



Different from conventional Fixed-Order phase-space generation (eg VEGAS)

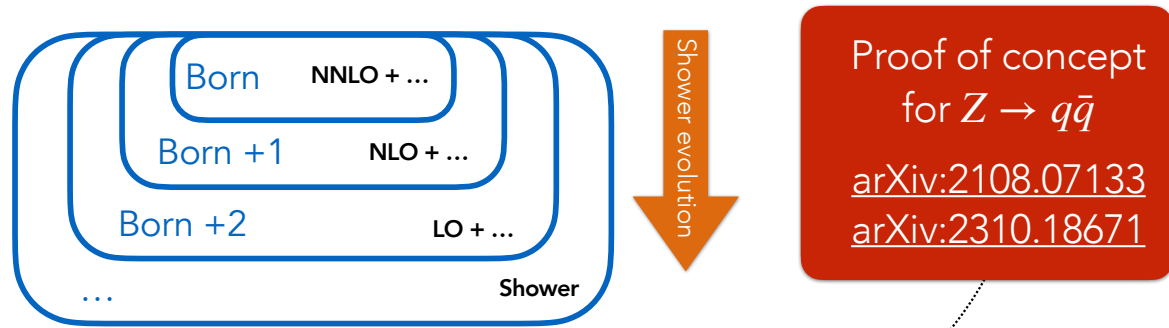


Towards True* NNLO Matching

*In the sense of the fixed-order and shower calculations matching each other point by point in each phase space

Continue shower afterwards

No auxiliary / unphysical scales \Rightarrow expect **small** matching systematics



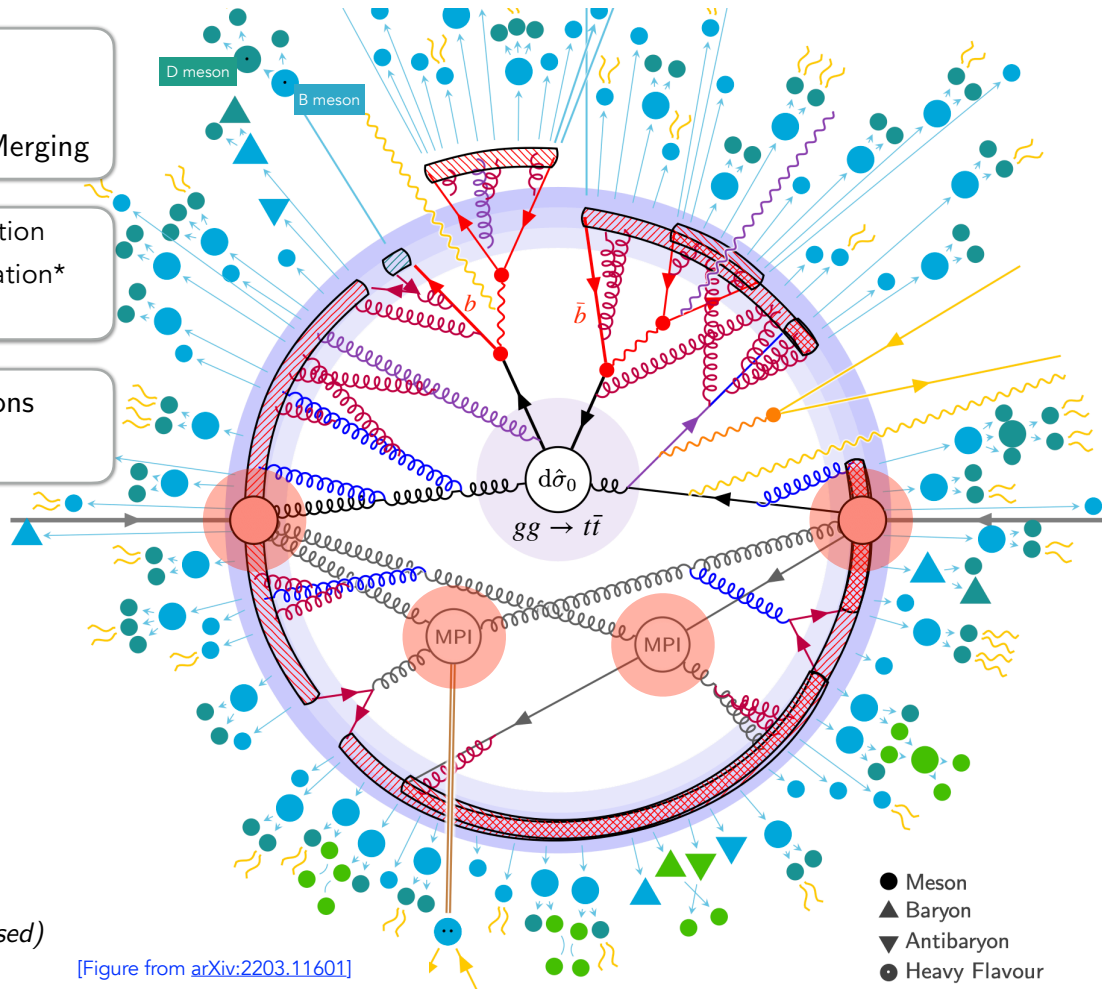
Need:

- 1 Born-Local NNLO ($\mathcal{O}(\alpha_s^2)$) K-factors: $k_{\text{NNLO}}(\Phi_2)$
- 2 NLO ($\mathcal{O}(\alpha_s^2)$) MECs in the first $2 \rightarrow 3$ shower emission: $k_{\text{NLO}}^{2 \rightarrow 3}(\Phi_3)$
- 3 LO ($\mathcal{O}(\alpha_s^2)$) MECs for next (iterated) $2 \rightarrow 3$ shower emission: $k_{\text{LO}}^{3 \rightarrow 4}(\Phi_4)$
- 4 Direct $2 \rightarrow 4$ branchings for unordered sector, with LO ($\mathcal{O}(\alpha_s^2)$) MECs: $k_{\text{LO}}^{2 \rightarrow 4}(\Phi_4)$



Part II – Nonperturbative Aspects

Hard Process	○ Hard Interaction
	● Resonance Decays
	■ MECs, Matching & Merging
Parton Showers	■ QCD Final-State Radiation
	■ QCD Initial-State Radiation*
	■ Electroweak Radiation
Underlying Event	○ Multiparton Interactions
	■ Beam Remnants*

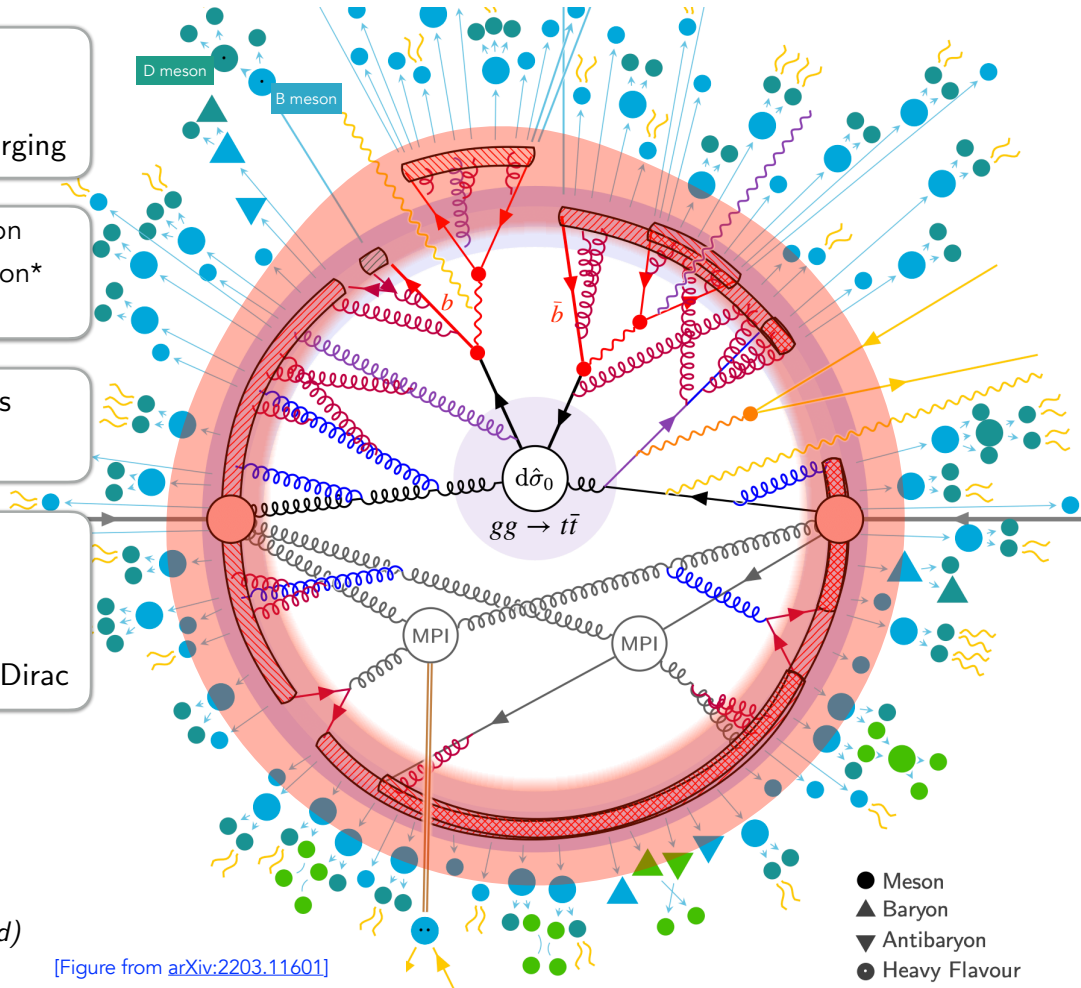
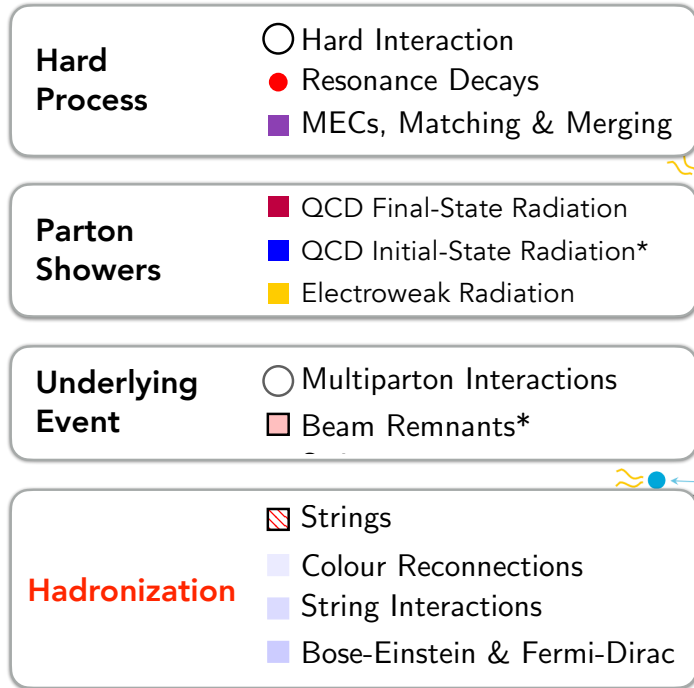


(*: incoming lines are crossed)

[Figure from [arXiv:2203.11601](https://arxiv.org/abs/2203.11601)]

- Meson
- ▲ Baryon
- ▼ Antibaryon
- Heavy Flavour

Hadronization



(*: incoming lines are crossed)

[Figure from [arXiv:2203.11601](https://arxiv.org/abs/2203.11601)]

- Meson
- ▲ Baryon
- ▼ Antibaryon
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Hadron Decays

Hard Process

- Hard Interaction
- Resonance Decays
- MECs, Matching & Merging

Parton Showers

- QCD Final-State Radiation
- QCD Initial-State Radiation*
- Electroweak Radiation

Underlying Event

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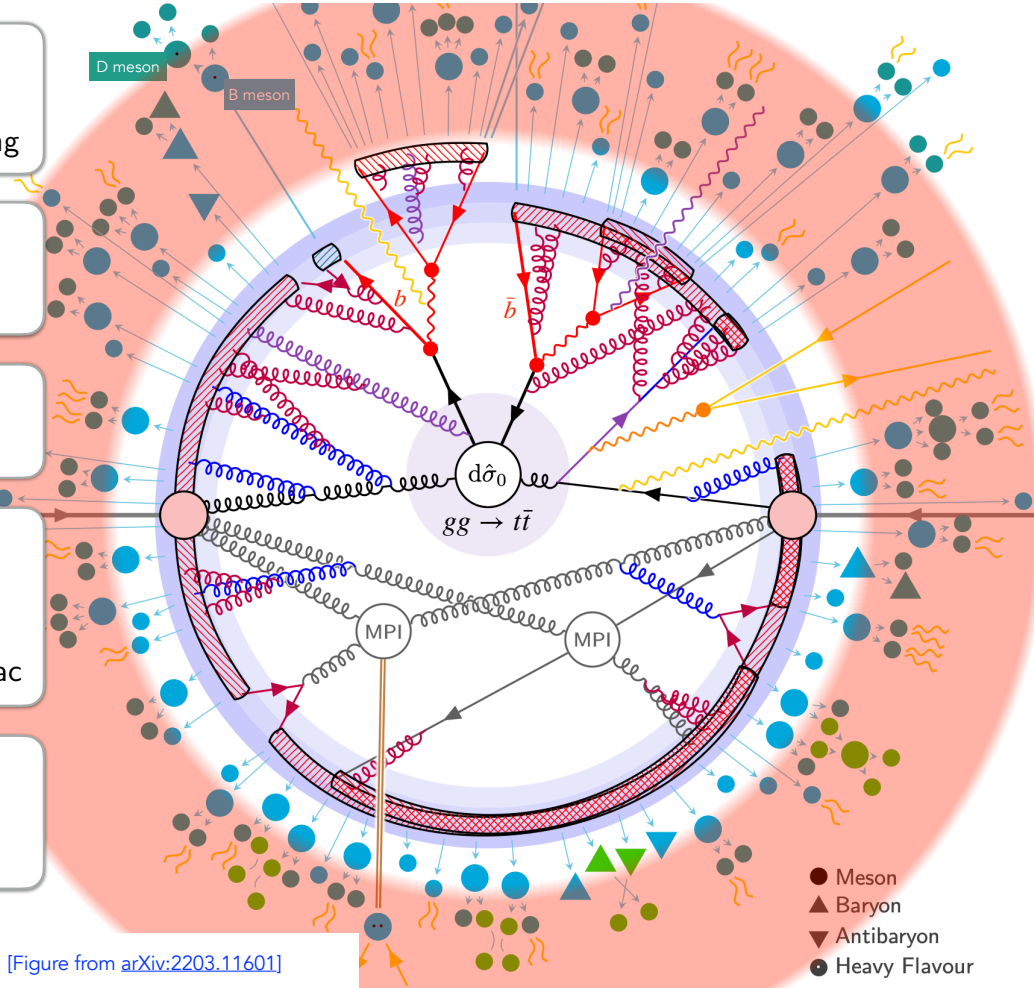
Hadronization

- ▨ Strings
- Colour Reconnections
- String Interactions
- Bose-Einstein & Fermi-Dirac

Hadron (& τ) Decays

- Primary Hadrons
- Secondary Hadrons
- Hadronic Reinteractions

(*: incoming lines are crossed)

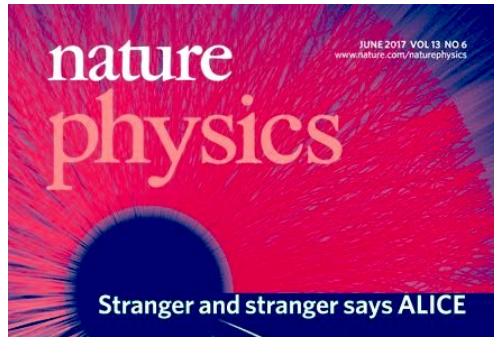


[Figure from arXiv:2203.11601]

New Discoveries in Hadronization

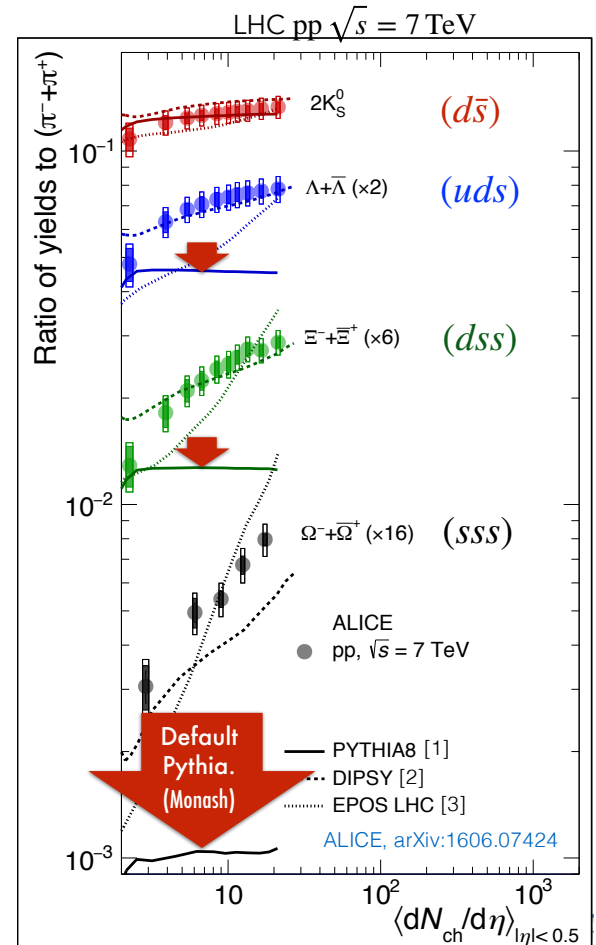
What a **strange** world we live in, said ALICE

Ratios of strange hadrons to pions strongly increase with event activity



June
2017

Conventional models (eg
Default PYTHIA) → **constant
strangeness fractions**



New Discoveries in Hadronization

LHC experiments also report very large (factor-10) enhancements in heavy-flavour baryon-to-meson ratios at low p_T !

Conventional models (eg default PYTHIA) \rightarrow constant baryon-to-meson ratio

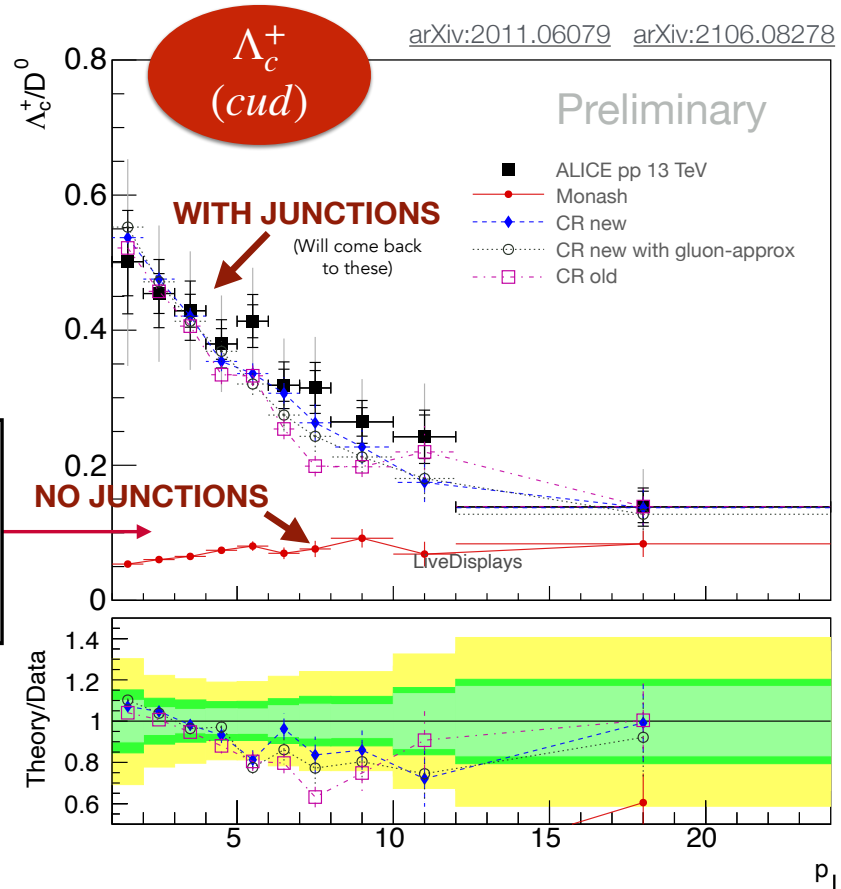
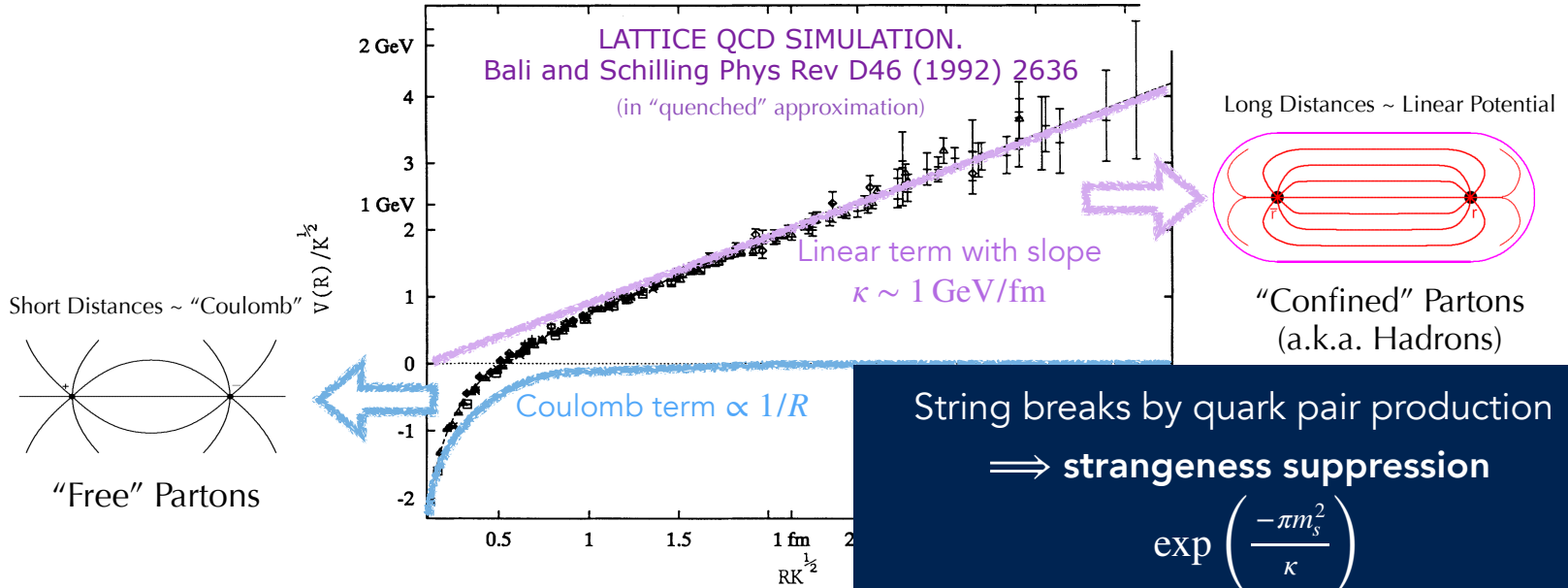


Figure from Altmann & PZS, *String Junctions Revisited*, in progress

Back to Basics – Anatomy of (Linear) Confinement

On lattice, compute potential energy of a colour-singlet $q\bar{q}$ state, as function of the distance, R , between the q and \bar{q} :



→ Model as strings (Lund Model)

String breaks by quark pair production
⇒ strangeness suppression

$$\propto \frac{\exp\left(\frac{-\pi m_s^2}{\kappa}\right)}{\exp\left(\frac{-\pi m_{u,d}^2}{\kappa}\right)}$$

Beyond the Static Limit

Regard tension κ as an emergent quantity?

Not fundamental strings

May depend on (invariant) time τ

E.g., hot strings which cool down

Hunt-Smith & **PZS** EPJC 80 (2020) 11

May depend on σ (excitations)

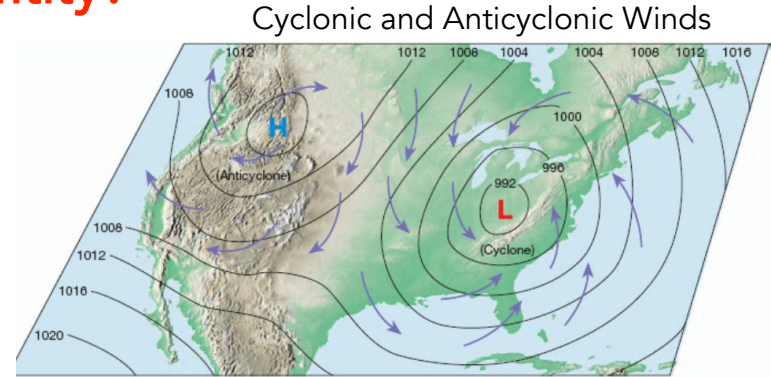
Working with E. Carragher & J. March-Russell in Oxford.

May depend on environment (e.g., other strings nearby)

Two approaches (so far) within Lund string-model context:

Colour Ropes [Bierlich, Gustafson, Lönnblad, Tarasov JHEP 03 (2015) 148; + more recent...]

Close-Packing [Fischer & Sjöstrand JHEP 01 (2017) 140; Altmann & **PZS** in progress ...]

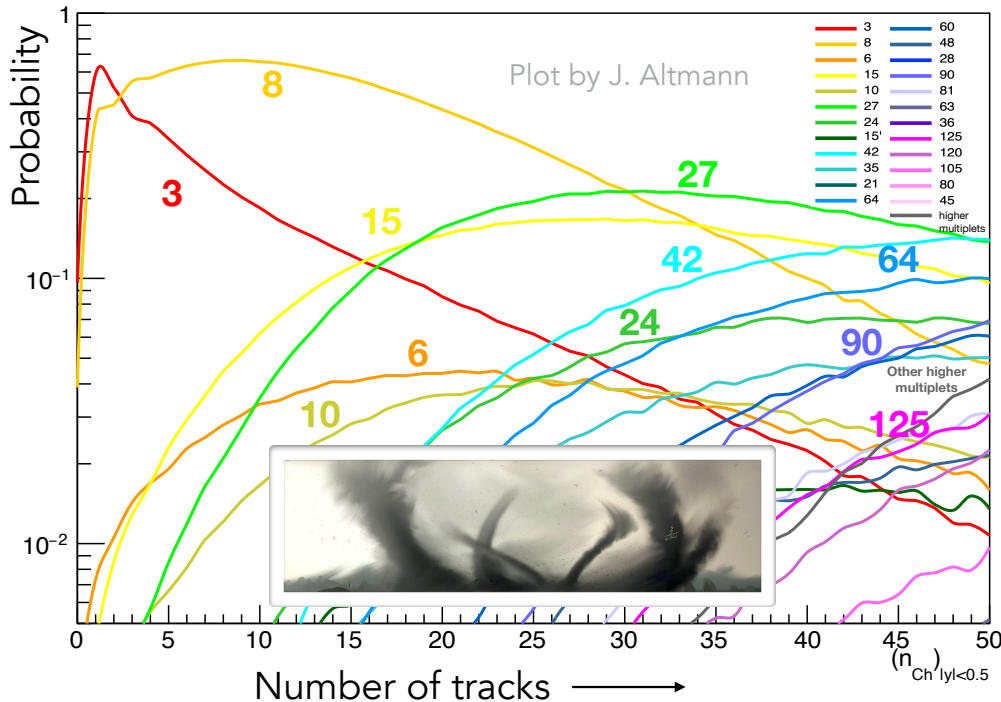


Non-Linear String Dynamics

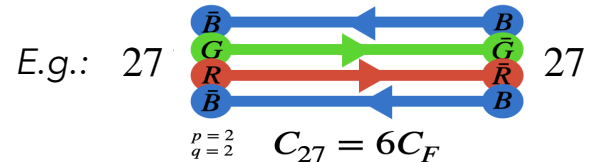
MPI \Rightarrow **lots** of coloured partons scattered into the final states

Count # of (oriented) flux lines crossing $y = 0$ in pp collisions (according to PYTHIA)

And classify by SU(3) multiplet:



Confining fields may be reaching **higher effective representations** than simple $q\bar{q}$ (3) ones.



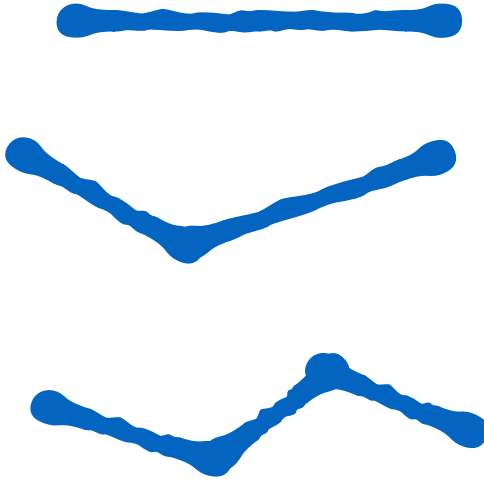
\rightarrow Is "emergent tension" driving strangeness enhancement in pp?

Altmann & PZS work in progress ...

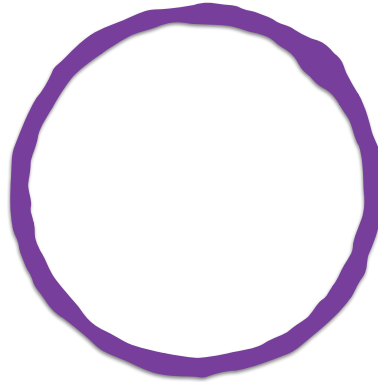
What about Baryon Number?

Types of string topologies:

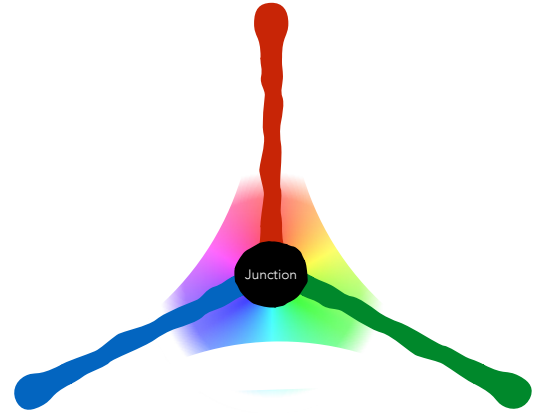
Open Strings



Closed Strings



SU(3) String Junction



Could we get these at LHC?

Stochastic sampling of $SU(3)$ group probabilities (e.g., $3 \otimes 3 = 6 \oplus \bar{3}$)

⇒ Random (re)connections in colour space (weighted by group weights)

Christiansen & PZS 2015

"QCD Colour Reconnections"

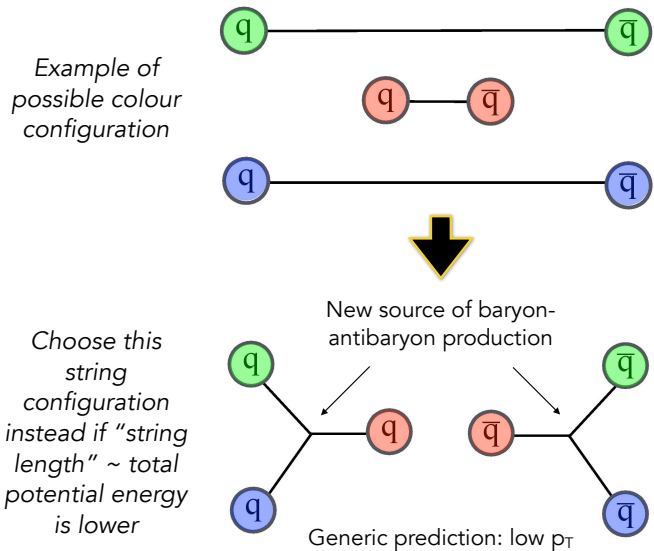
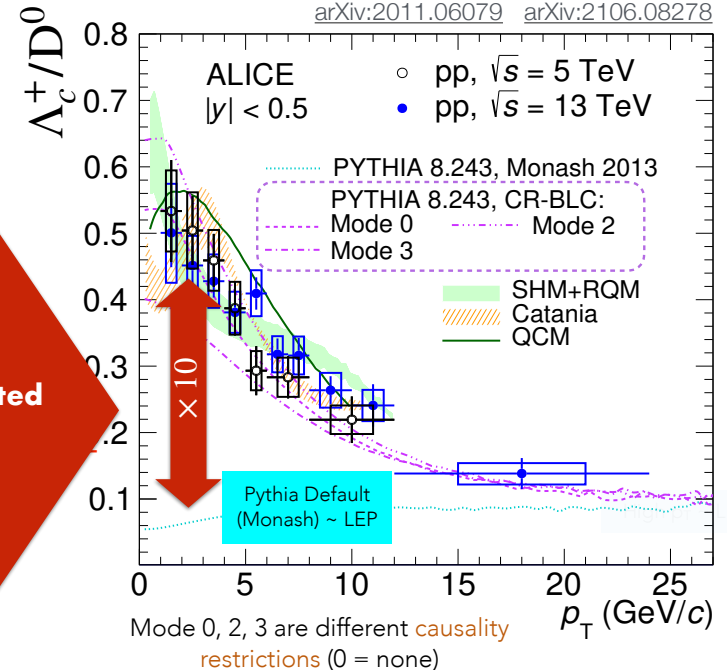


Illustration by J. Altmann

ALICE 2021

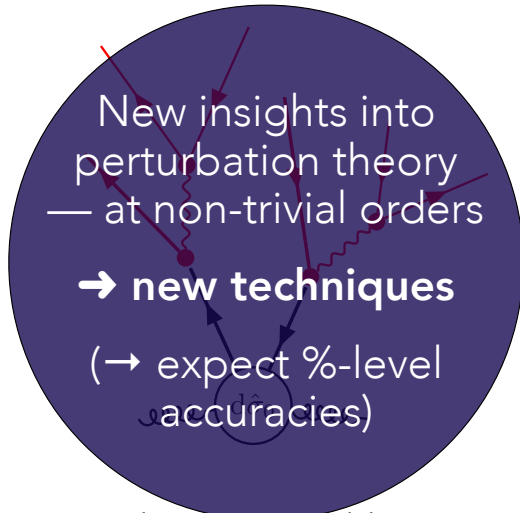


Summary

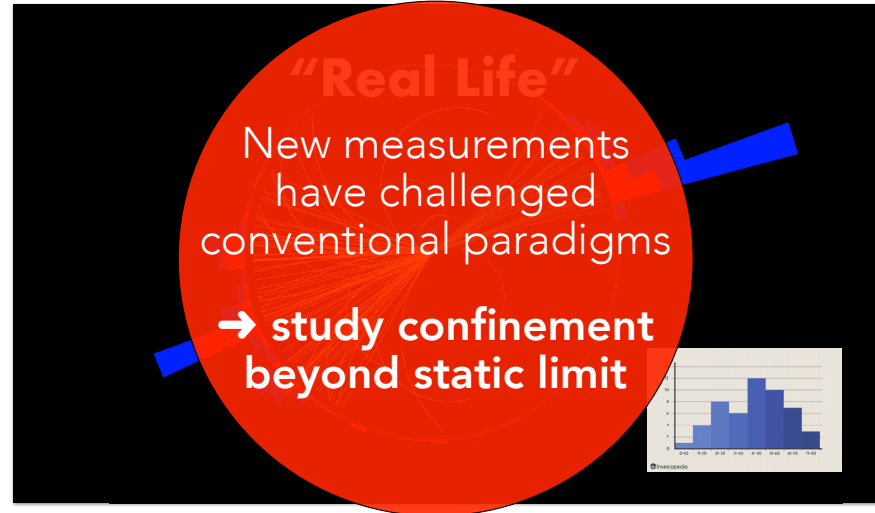
The Goal: use LHC measurements to test hypotheses about Nature

Problem 1: no **exact** solutions to QFT

→ Perturbative **Approximations**



Elementary Fields,
Symmetries,
Interactions



Problem 2: Confinement

We collide — and observe — **hadrons**