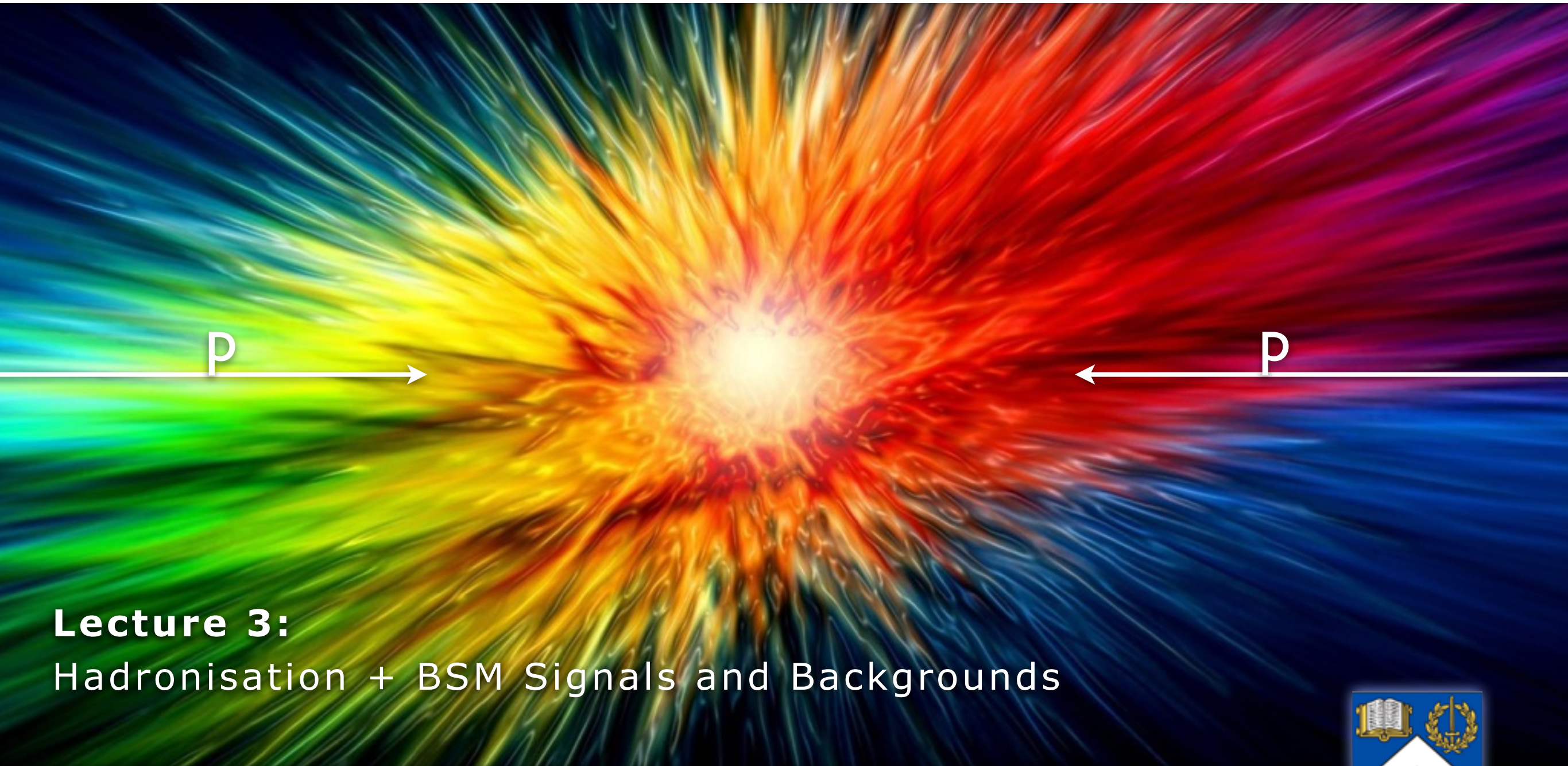


Monte Carlos and New Physics

Peter Skands (Monash University)



Lecture 3:

Hadronisation + BSM Signals and Backgrounds

Pre-SUSY - June 2016

Lecture Notes: [P. Skands, arXiv:1207.2389](https://arxiv.org/abs/1207.2389)



Monte Carlos & Fragmentation

PYTHIA anno 1978 (then called JETSET)

LU TP 78-18
November, 1978

A Monte Carlo Program for Quark Jet
Generation

T. Sjöstrand, B. Söderberg

A Monte Carlo computer program is presented, that simulates the **fragmentation of a fast parton into a jet of mesons**. It uses an iterative scaling scheme and is compatible with the jet model of Field and Feynman.

Note:

Field-Feynman was an early fragmentation model
Now superseded by the String (in PYTHIA) and
Cluster (in HERWIG & SHERPA) models.

```
SUBROUTINE JETGEN(N)
COMMON /JET/ K(100,2), P(100,5)
COMMON /PAR/ PUD, PS1, SIGMA, CX2, EBEG, WFIN, IFLBEG
COMMON /DATA1/ MESO(9,2), CMIX(6,2), PMAS(19)
IFLSGN=(10-IFLBEG)/5
W=2.*EBEG
I=0
IPD=0
C 1 FLAVOUR AND PT FOR FIRST QUARK
IFL1=IABS(IFLBEG)
PT1=SIGMA*SQRT(-ALOG(RANF(0)))
PHI1=6.2832*RANF(0)
PX1=PT1*COS(PHI1)
PY1=PT1*SIN(PHI1)
100 I=I+1
C 2 FLAVOUR AND PT FOR NEXT ANTIQUARK
IFL2=1+INT(RANF(0)/PUD)
PT2=SIGMA*SQRT(-ALOG(RANF(0)))
PHI2=6.2832*RANF(0)
PX2=PT2*COS(PHI2)
PY2=PT2*SIN(PHI2)
C 3 MESON FORMED, SPIN ADDED AND FLAVOUR MIXED
K(I,1)=MESO(3*(IFL1-1)+IFL2,IFLSGN)
ISPIN=INT(PS1+RANF(0))
K(I,2)=1+9*ISPIN+K(I,1)
IF(K(I,1).LE.6) GOTO 110
TMIX=RANF(0)
KM=K(I,1)-6+3*ISPIN
K(I,2)=8+9*ISPIN+INT(TMIX+CMIX(KM,1))+INT(TMIX+CMIX(KM,2))
C 4 MESON MASS FROM TABLE, PT FROM CONSTITUENTS
110 P(I,5)=PMAS(K(I,2))
P(I,1)=PX1+PX2
P(I,2)=PY1+PY2
PMTS=P(I,1)**2+P(I,2)**2+P(I,5)**2
C 5 RANDOM CHOICE OF X=(E+PZ)MESON/(E+PZ)AVAILABLE GIVES E AND PZ
X=RANF(0)
IF(RANF(0).LT.CX2) X=1.-X**(1./3.)
P(I,3)=(X*W-PMTS/(X*W))/2.
P(I,4)=(X*W+PMTS/(X*W))/2.
C 6 IF UNSTABLE, DECAY CHAIN INTO STABLE PARTICLES
120 IPD=IPD+1
IF(K(IPD,2).GE.8) CALL DECAY(IPD,I)
IF(IPD.LT.1.AND.I.LE.96) GOTO 120
C 7 FLAVOUR AND PT OF QUARK FORMED IN PAIR WITH ANTIQUARK ABOVE
IFL1=IFL2
PX1=-PX2
PY1=-PY2
C 8 IF ENOUGH E+PZ LEFT, GO TO 2
W=(1.-X)*W
IF(W.GT.WFIN.AND.I.LE.95) GOTO 100
N=I
RETURN
END
```

Fast-Forward ~ 40 Years

PYTHIA anno 2016

(now called PYTHIA 8)

~ 100,000 lines of C++

What a modern MC generator has inside:

CPC 191 (2015) p.159-177
October, 2014

An Introduction to PYTHIA 8.2

T. Sjöstrand et al. (10 authors)

The Pythia program is a standard tool for the generation of events in high-energy collisions, comprising a coherent set of physics models for the evolution from a few-body hard process to a complex multiparticle final state. It contains a library of hard processes, models for initial- and final-state parton showers, matching and merging methods between hard processes and parton showers, multiparton interactions, beam remnants, string fragmentation and particle decays. It also has a set of utilities and several interfaces ...

- Hard Processes (internal, interfaced, or via Les Houches events)
- BSM (internal or via interfaces)
- PDFs (internal or via interfaces)
- Showers (internal or inherited)
- Multiple parton interactions
- Beam Remnants
- String Fragmentation
- Decays (internal or via interfaces)
- Examples and Tutorial
- Online HTML / PHP Manual
- Utilities and interfaces to external programs

The Main Workhorses



PYTHIA (begun 1978)

Originated in hadronisation studies: Lund String model
Still special emphasis on soft physics



HERWIG (begun 1984)

Originated in coherence studies: angular-ordered showers
Cluster hadronisation as simple complement



SHERPA (begun ~2000)

Originated in “matching” of matrix elements to showers (CKKW-L), with own internal matrix-element generator(s)

+ Many more specialised:

Matrix-Element Generators, Matching/Merging Packages

Soft-QCD, Cosmic-Ray, and Heavy-Ion Generators

(BSM) Model Generators, Decay Packages, Alternative QCD showers, ...

From Partons to Pions

Here's a fast parton

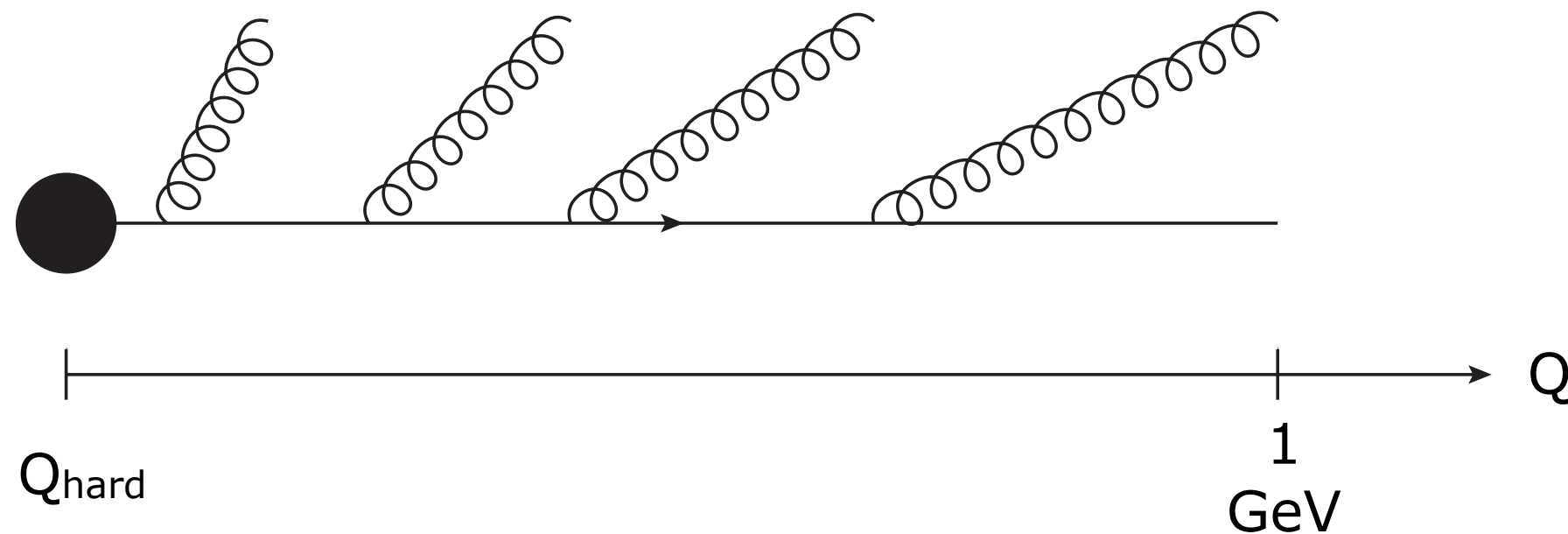
Fast: It starts at a high factorization scale

$$Q = Q_F = Q_{\text{hard}}$$

It showers
(bremsstrahlung)

It ends up
at a low effective
factorization scale

$$Q \sim m_\rho \sim 1 \text{ GeV}$$



From Partons to Pions

Here's a fast parton

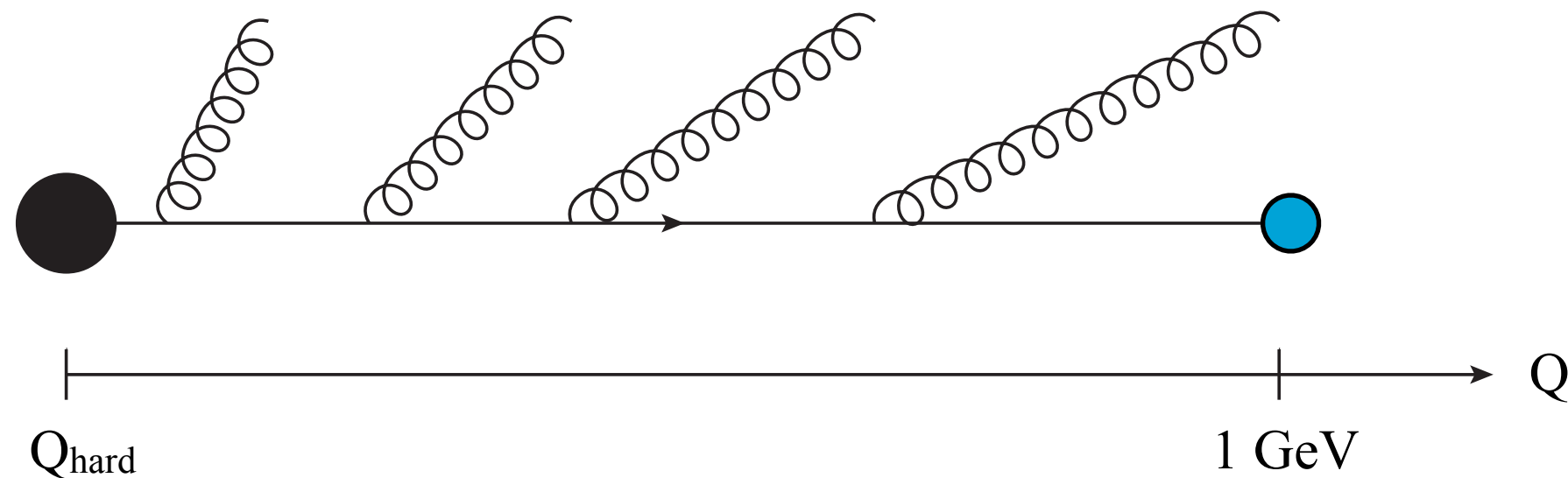
Fast: It starts at a high factorization scale

$$Q = Q_F = Q_{\text{hard}}$$

It showers
(bremsstrahlung)

It ends up
at a low effective
factorization scale

$$Q \sim m_p \sim 1 \text{ GeV}$$



How about I just call it a hadron?

→ "Local Parton-Hadron Duality"

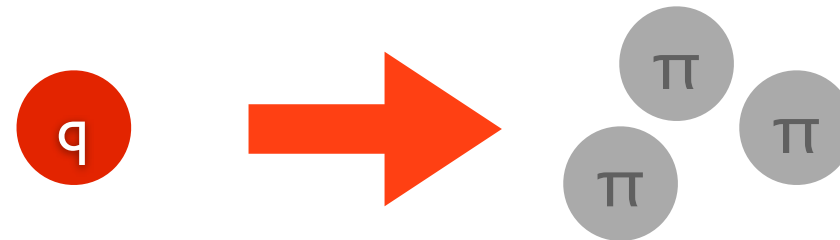
Parton \rightarrow Hadrons?

Early models: “Independent Fragmentation”

Local Parton Hadron Duality (LPHD) can give useful results for inclusive quantities in collinear fragmentation

Motivates a simple model:

“Independent Fragmentation”



But ...

The point of confinement is that partons are coloured

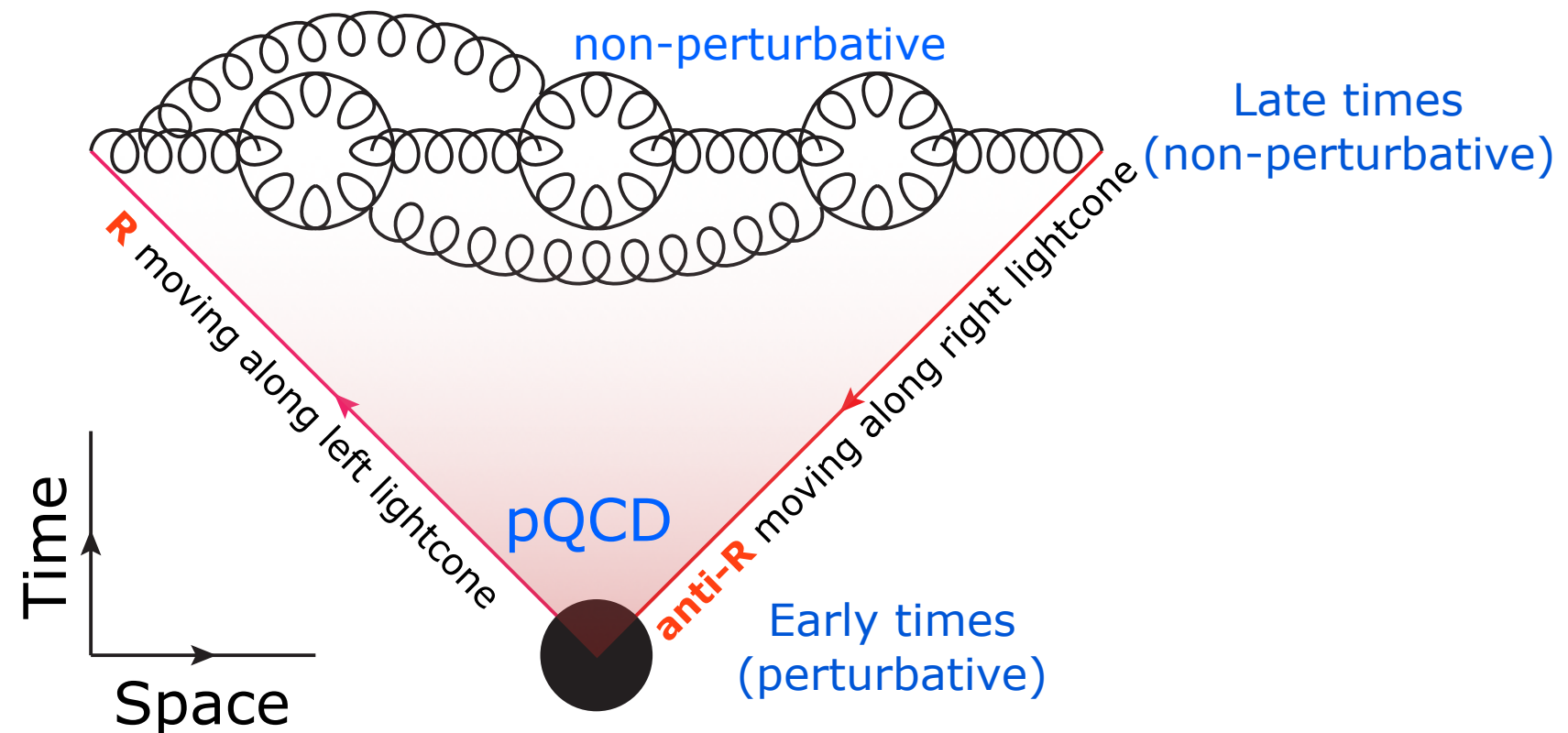
Hadronisation = the process of **colour neutralisation**

- Unphysical to think about independent fragmentation of a single parton into hadrons
- Too naive to see LPHD (inclusive) as a justification for Independent Fragmentation (exclusive)
- More physics needed

Colour Neutralisation

A physical hadronization model

Should involve at least TWO partons, with opposite color charges (e.g., **R** and **anti-R**)



Strong “confining” field emerges between the two charges when their separation $> \sim 1\text{fm}$

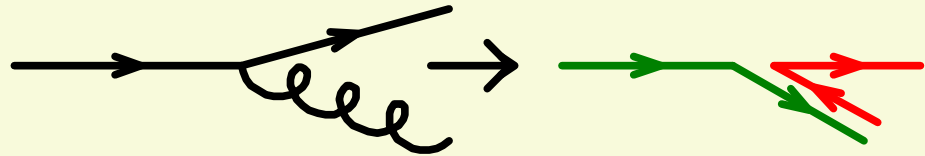
Which Charges? Colour Flow

After the parton shower finishes, there can be lots of partons, $\mathcal{O}(10-100)$. The main question is therefore:

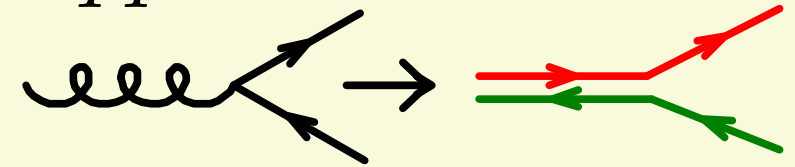
Between which partons do confining potentials arise?

MC generators use a simple set of rules for color flow, based on large- N_c limit (valid to $\sim 1/N_c^2 \sim 10\%$)

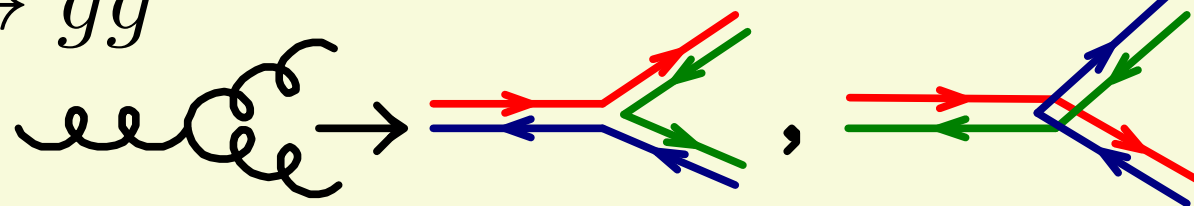
$q \rightarrow qg$



$g \rightarrow q\bar{q}$



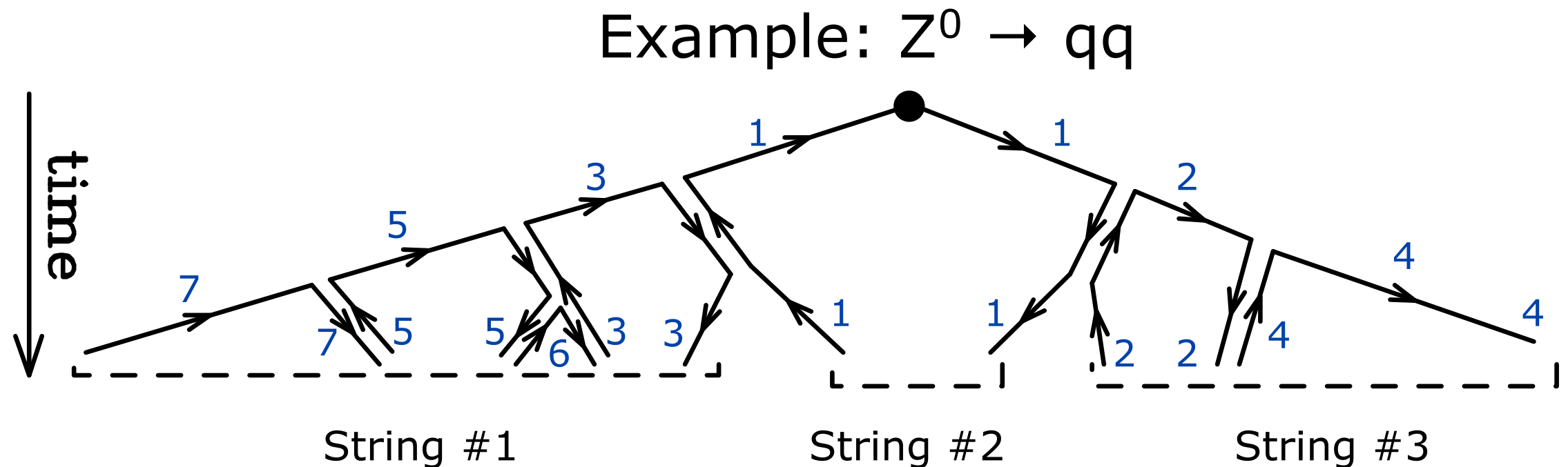
$g \rightarrow gg$



Illustrations from: Nason & Skands, PDG Review on *MC Event Generators*, 2014

Color Flow

For an entire Cascade



Coherence of pQCD cascades \rightarrow not much "overlap" between strings
 \rightarrow Leading-colour approximation pretty good

(LEP measurements in $e^+e^- \rightarrow W^+W^- \rightarrow \text{hadrons}$ confirm this (at least to order 10% $\sim 1/N_c^2$))

Note: (much) more color getting kicked around in hadron collisions.

Signs that LC approximation is breaking down?

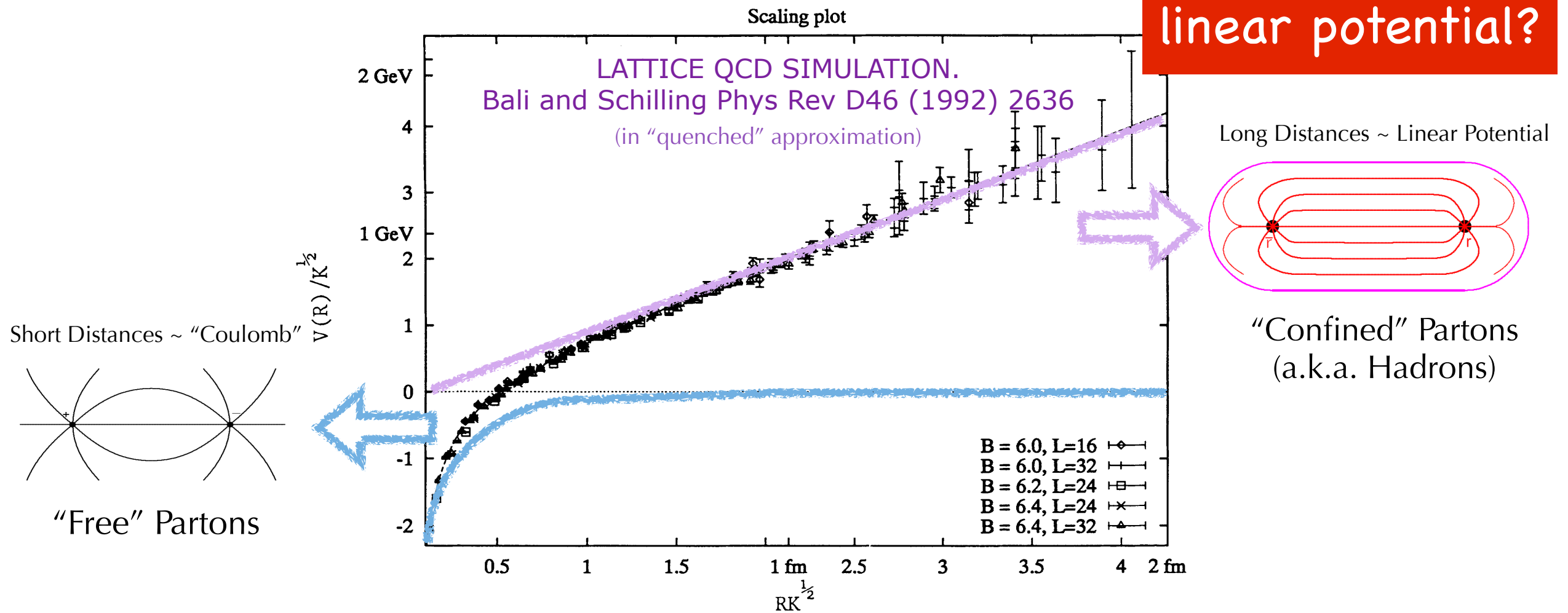
see, e.g., [Christiansen, Skands JHEP 1508 \(2015\) 003](#)

The Ultimate Limit: Wavelengths $> 10^{-15}$ m

Quark-Antiquark Potential

As function of separation distance

What physical system has a linear potential?



$$F(r) \approx \text{const} = \kappa \approx 1 \text{ GeV/fm} \iff V(r) \approx \kappa r$$

~ Force required to lift a 16-ton truck

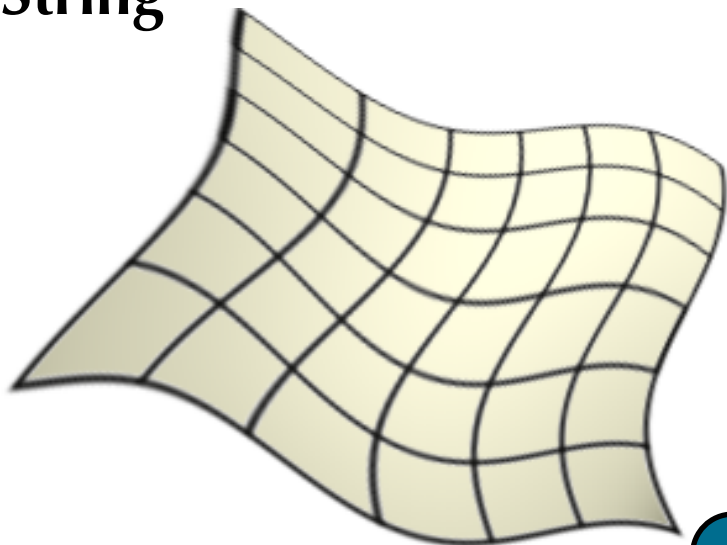
From Partons to Strings

Motivates a model:

Let color field collapse into a (infinitely) narrow flux tube of uniform energy density $\kappa \sim 1 \text{ GeV / fm}$

→ Relativistic 1+1 dimensional worldsheet

String



Pedagogical Review: B. Andersson, *The Lund model*. Camb. Monogr. Part. Phys. Nucl. Phys. Cosmol., 1997.

In “unquenched” QCD

$g \rightarrow qq \rightarrow$ The strings will break

Schwinger Effect

Non-perturbative creation of e^+e^- pairs in a strong external Electric field

Probability from Tunneling Factor

$$\mathcal{P} \propto \exp\left(\frac{-m^2 - p_{\perp}^2}{\kappa/\pi}\right)$$

(κ is the string tension equivalent)

→ Gaussian p_T spectrum
Heavier quarks suppressed. Prob($q=d,u,s,c$) $\approx 1 : 1 : 0.2 : 10^{-11}$

(Note on the Length of Strings)

In Space:

String tension ≈ 1 GeV/fm \rightarrow a 5-GeV quark can travel 5 fm before all its kinetic energy is transformed to potential energy in the string.

Then it must start moving the other way. String breaks will have happened behind it \rightarrow yo-yo model of mesons

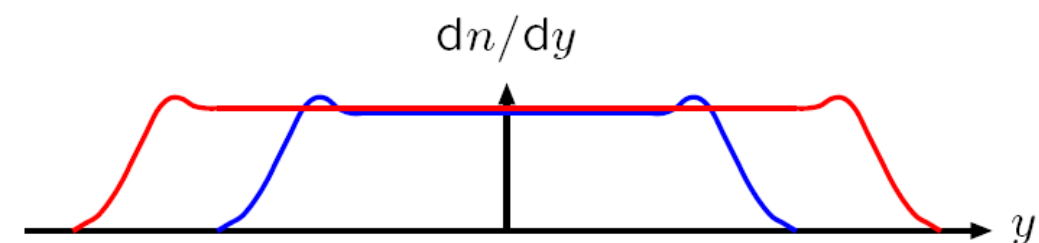
In Rapidity :
$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) = \frac{1}{2} \ln \left(\frac{(E + p_z)^2}{E^2 - p_z^2} \right)$$

For a pion with $z=1$ along string direction
(For beam remnants, use a proton mass):

$$y_{\max} \sim \ln \left(\frac{2E_q}{m_\pi} \right)$$

Note: Constant average hadron multiplicity per unit $y \rightarrow$ logarithmic growth of total multiplicity

Scaling in lightcone $p_\pm = E \pm p_z$ (for $q\bar{q}$ system along z axis) implies flat central rapidity plateau + some endpoint effects:

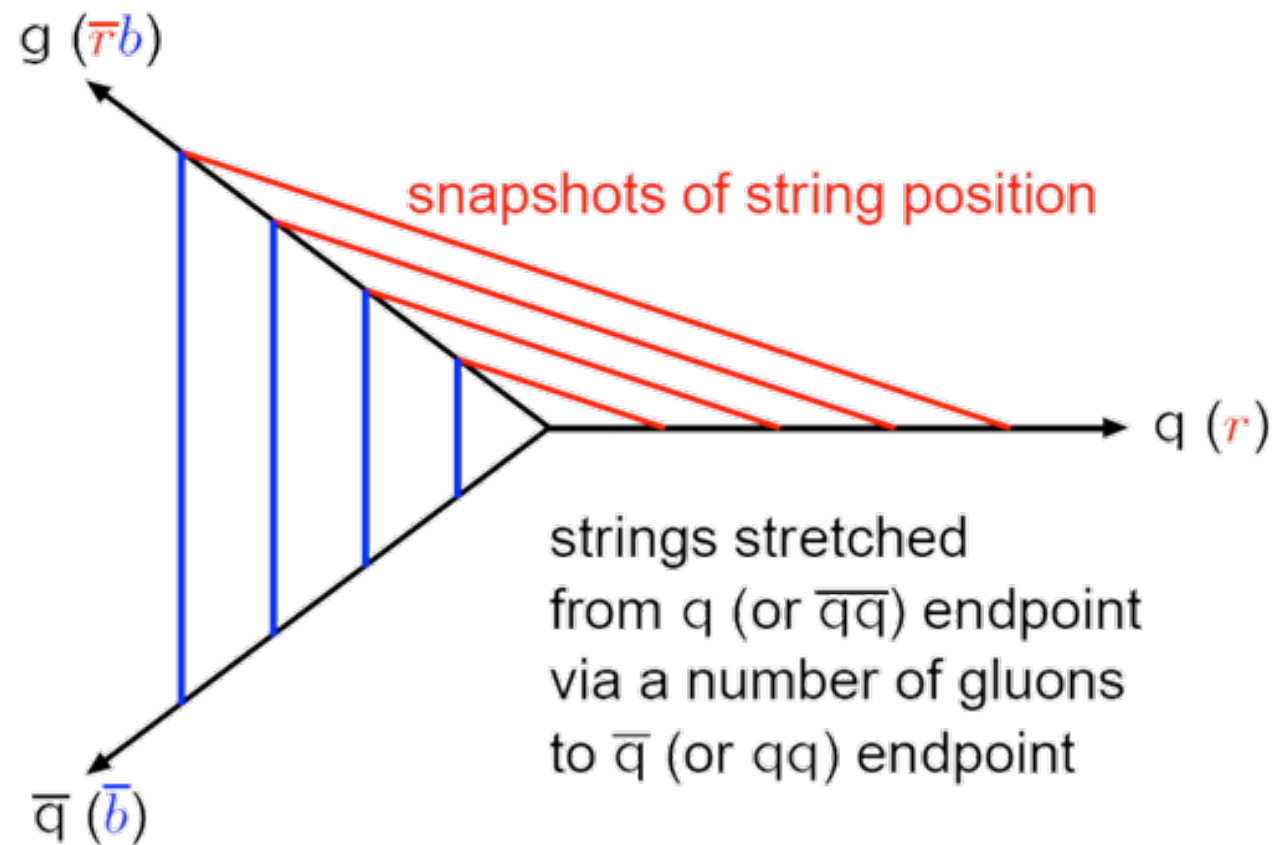


$$\langle n_{\text{ch}} \rangle \approx c_0 + c_1 \ln E_{\text{cm}}, \sim \text{Poissonian multiplicity distribution}$$

The (Lund) String Model

Map:

- **Quarks** → String Endpoints
- **Gluons** → Transverse Excitations (kinks)
- Physics then in terms of string worldsheet evolving in spacetime
- Probability of string break (by quantum tunneling) constant per unit area → **AREA LAW**



strings stretched from q (or \bar{q}) endpoint via a number of gluons to \bar{q} (or q) endpoint

Gluon = kink on string, carrying energy and momentum

→ **STRING EFFECT**

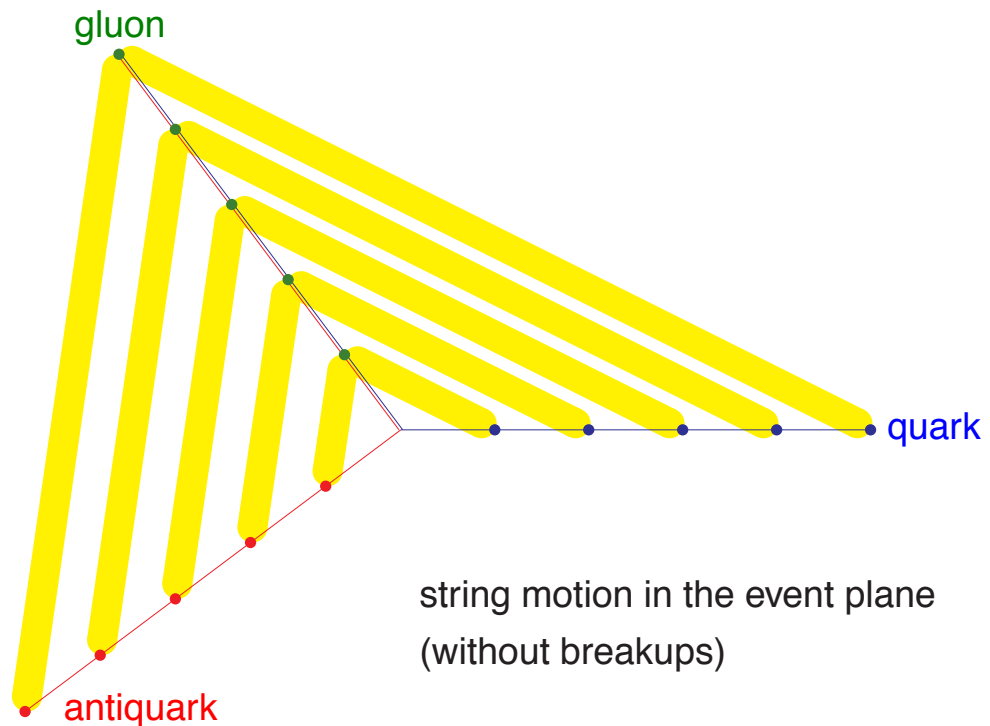
Simple space-time picture

Details of string breaks more complicated (e.g., baryons, spin multiplets)

Differences Between Quark and Gluon Jets

More recent study (LHC)

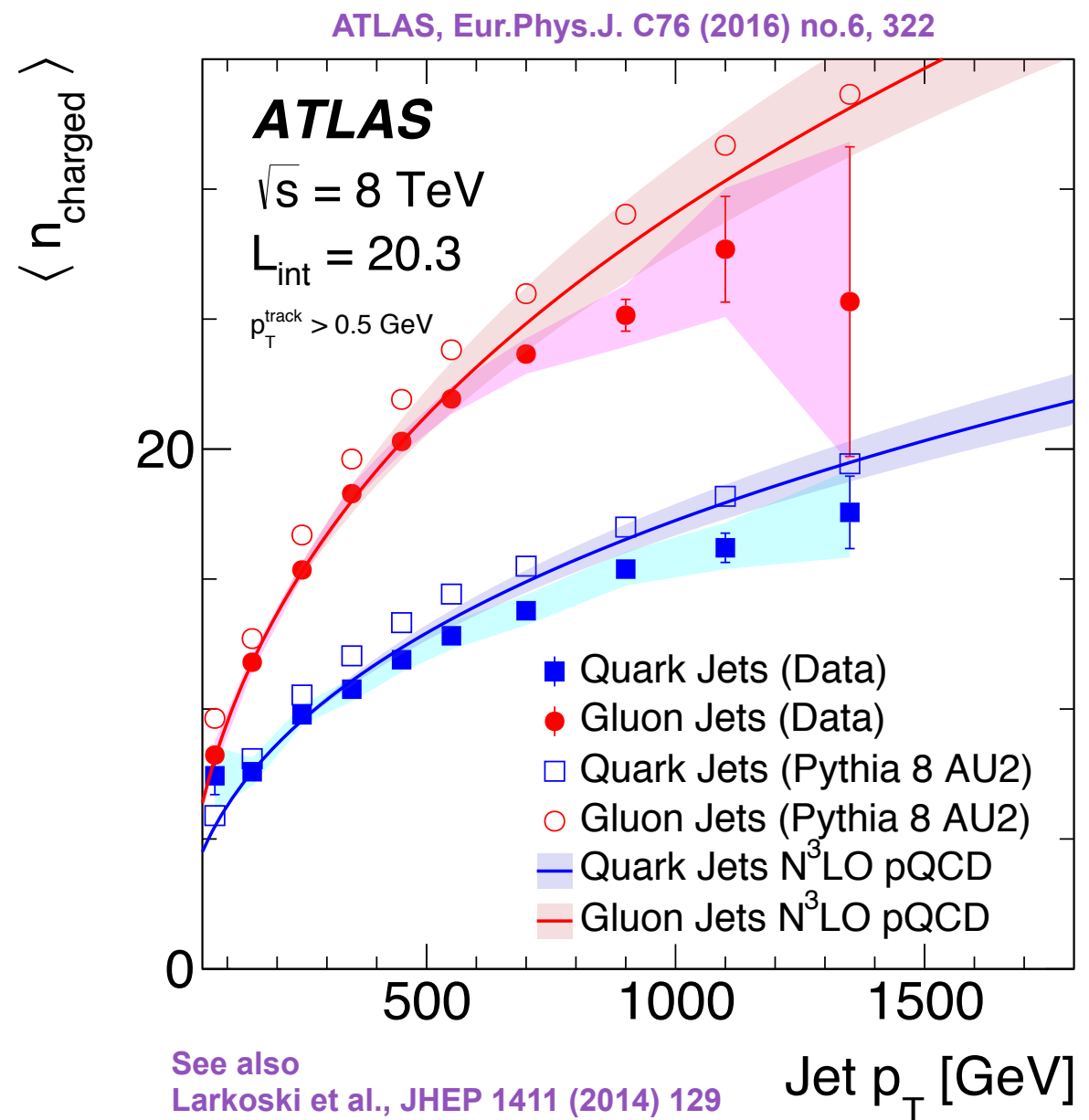
Gluon connected to two string pieces



Each quark connected to one string piece

→ expect factor $2 \sim C_A/C_F$ larger particle multiplicity in gluon jets vs quark jets

Can be hugely important for discriminating new-physics signals (decays to quarks vs decays to gluons, vs composition of background and bremsstrahlung combinatorics)



See also

Larkoski et al., JHEP 1411 (2014) 129

Thaler et al., Les Houches, arXiv:1605.04692

The Effects of Hadronisation

Generally, expect few-hundred MeV shifts by hadronisation

Corrections to IR safe observables are “power corrections”

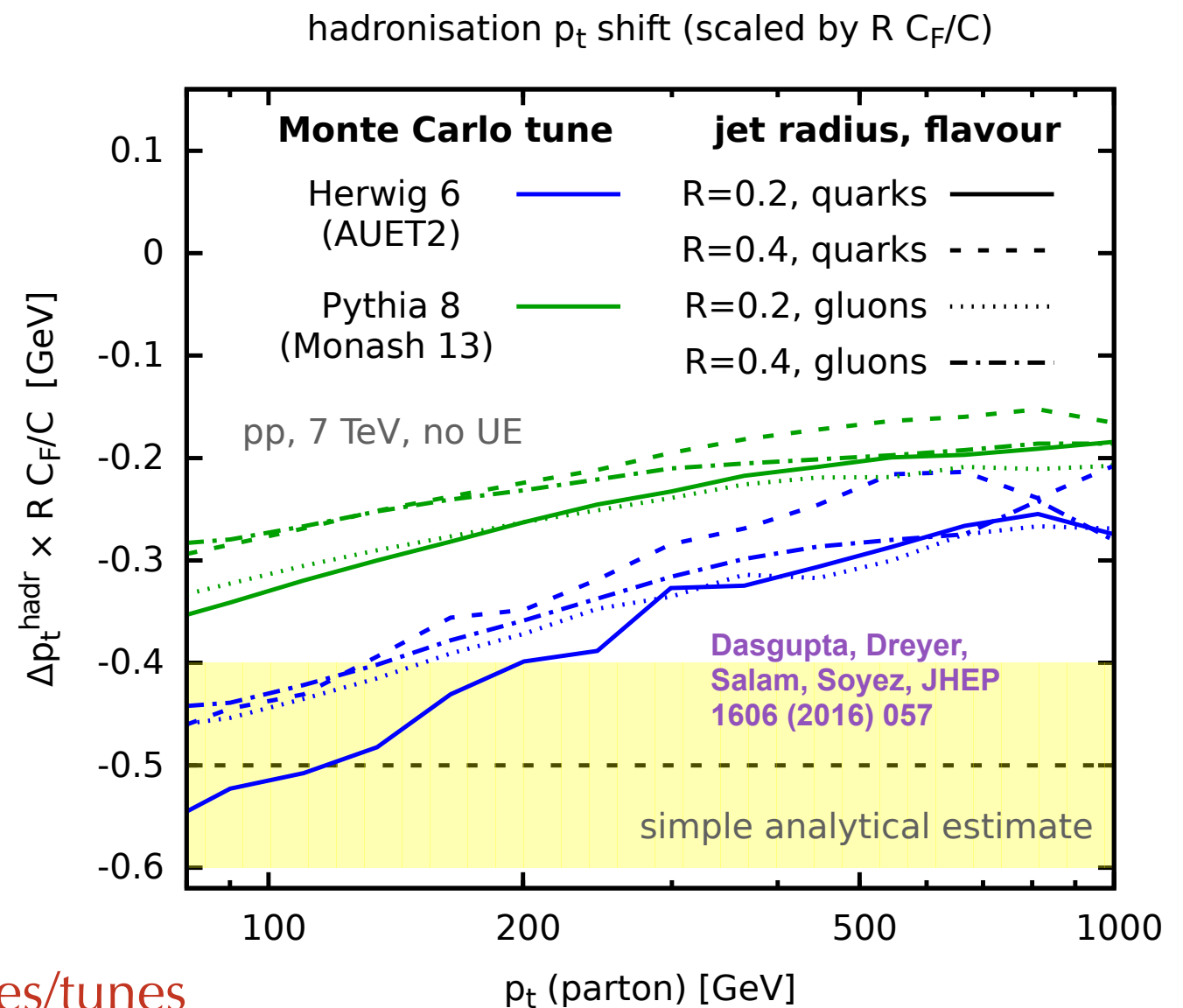
$$\propto \Lambda_{\text{QCD}}^2 / Q_{\text{OBS}}^2$$

Corrections for jets
of radius $R = \Delta\eta \times \Delta\phi$

$$\propto 1/R$$

See
Korchensky, Sterman, NPB 437 (1995) 415
Seymour, NPB 513 (1998) 269
Dasgupta, Magnea, Salam, JHEP 0802 (2008) 055

Simple analytical estimate
→ $\sim 0.5 \text{ GeV} / R$ correction
from hadronisation
(scaled by colour factor)



Significant differences between codes/tunes

→ important to pin down with precise QCD hadronisation measurements at LHC

(Alternative: The Cluster Model)

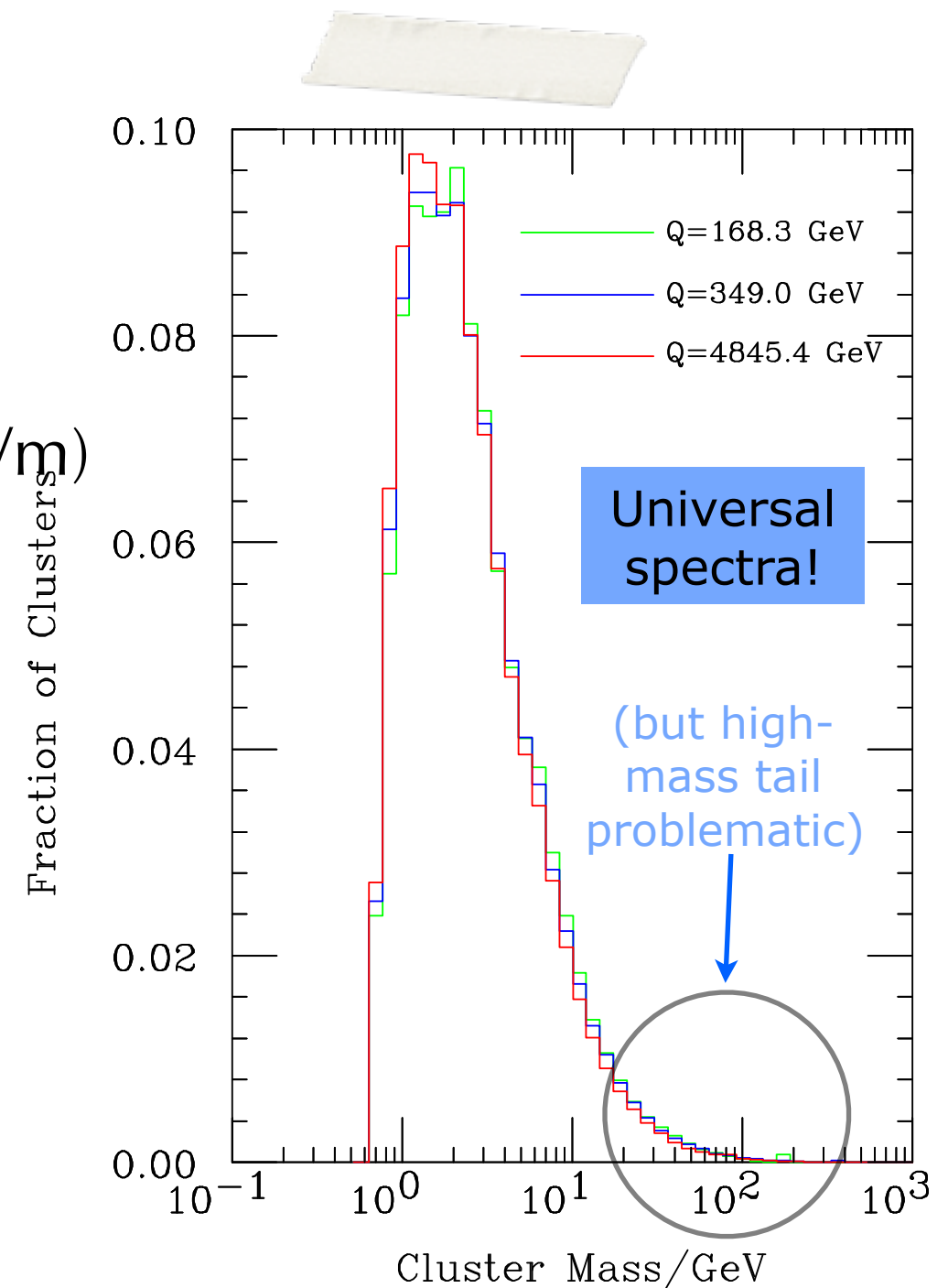
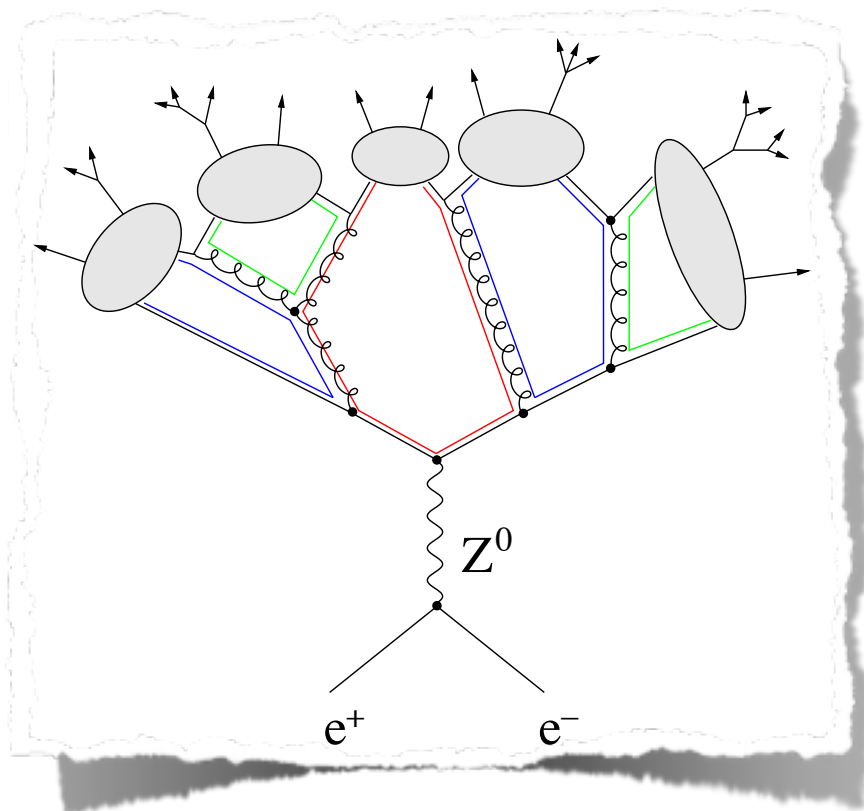
“Preconfinement”

+ Force $g \rightarrow qq$ splittings at Q_0

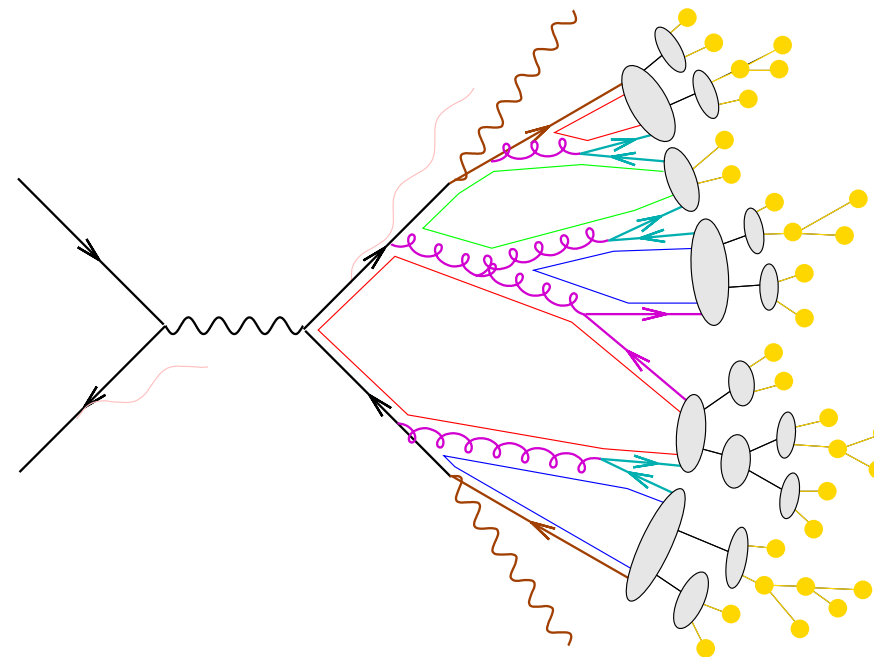
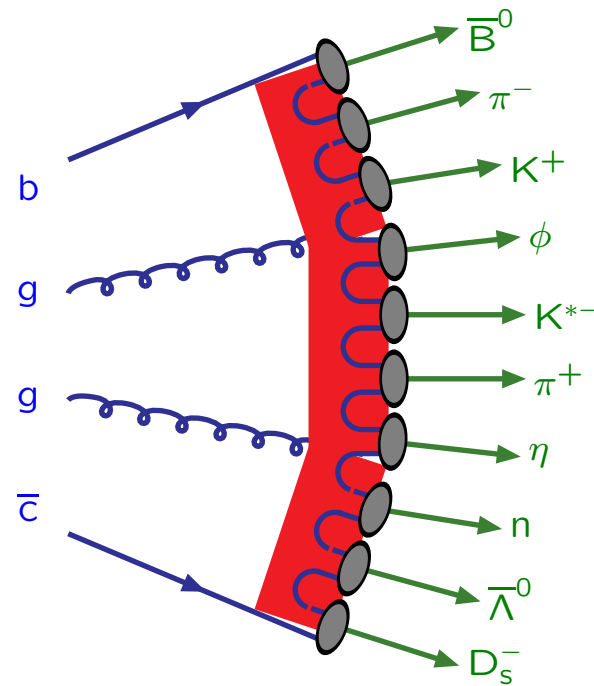
→ high-mass q - q bar “clusters”

Isotropic 2-body decays to hadrons

according to PS $\approx (2s_1+1)(2s_2+1)(p^*/m)$



Strings and Clusters



program model	PYTHIA string	HERWIG & SHERPA cluster
energy-momentum picture	powerful	simple
parameters	predictive	unpredictive
flavour composition	few	many
parameters	messy	simple
	unpredictive	in-between
	many	few

Small strings → clusters. Large clusters → strings

Monte Carlos and New Physics

Aspects where MC is needed (apart from background modelling)

Signal properties

Decay distributions, extra jets, jet structure (\rightarrow jet calibrations),
QCD (& EW) corrections to kinematic distributions

Exclusions & “Recasting” of searches

Uncertainties / modelling deficiencies \rightarrow worse exclusions

Dynamical modelling of BSM phenomena (some examples)

Long-Lived Coloured Particles \rightarrow “R-hadrons”

Dark-Matter Annihilation to Coloured Particles

Hidden Valleys (showers/hadronisation in hidden sector)

Baryon Number Violation (RPV-SUSY) \rightarrow colour-epsilon structures

Particles with “exotic” colour charges (e.g., colour sextets)

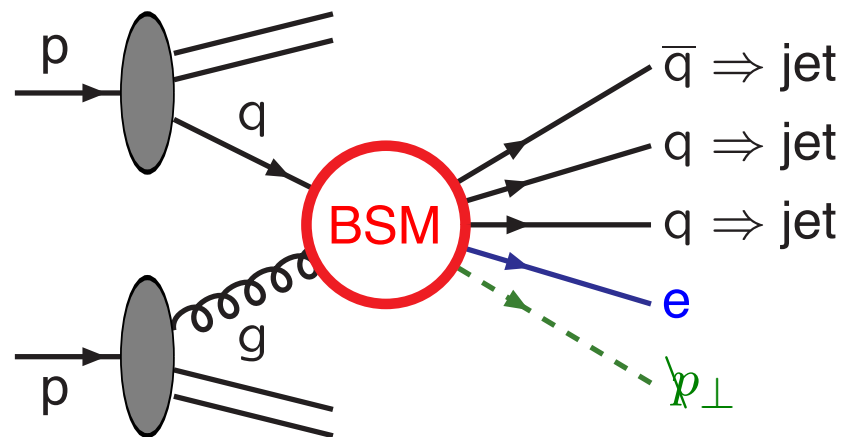
Black-Hole Evaporation, Sphaleron Decays, ...

BSM at the LHC

(adapted from slides by T. Sjostrand)

BSM particles usually short-lived, or weakly interacting (like DM)

Then visible final state consists of hadrons, leptons and photons, just like ordinary processes.



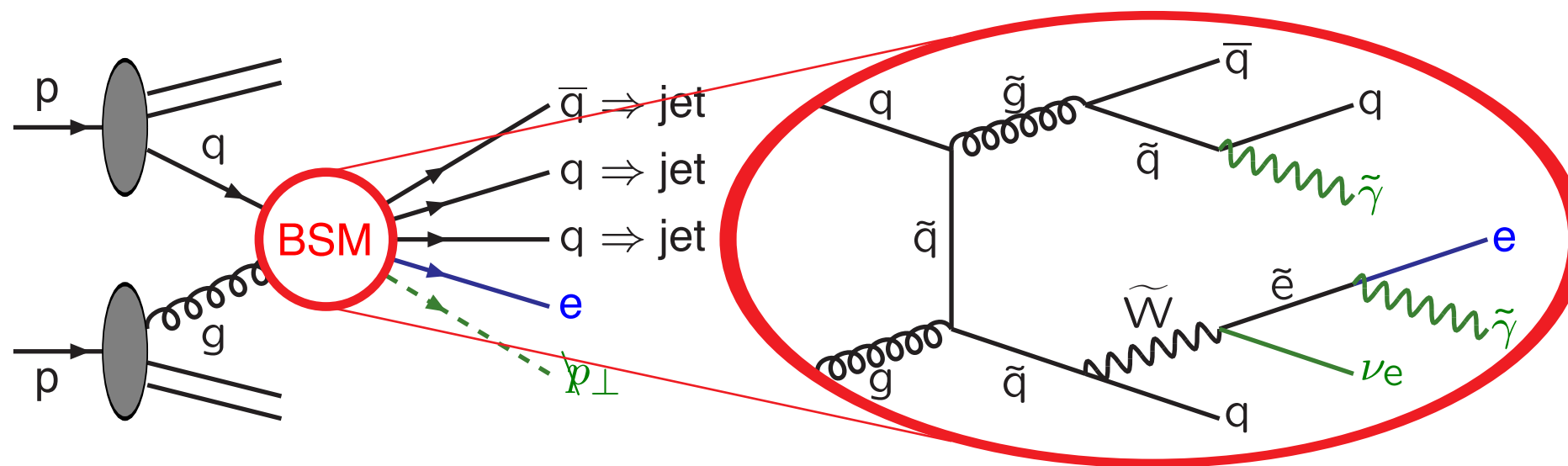
As easy to model as SM processes.

BSM at the LHC

(adapted from slides by T. Sjostrand)

BSM particles usually short-lived, or weakly interacting (like DM)

Then visible final state consists of hadrons, leptons and photons, just like ordinary processes.



As easy to model as SM processes.

Original structure hidden, but traces of it may be left in terms of invariant masses and angular distributions

Discovery requires detailed understanding of rare signals and huge backgrounds

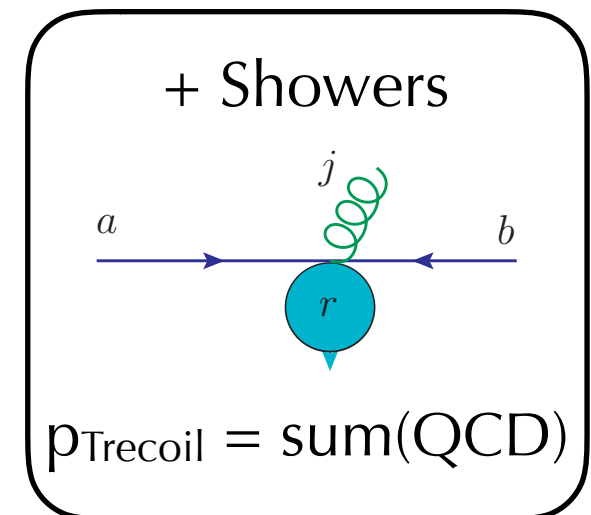
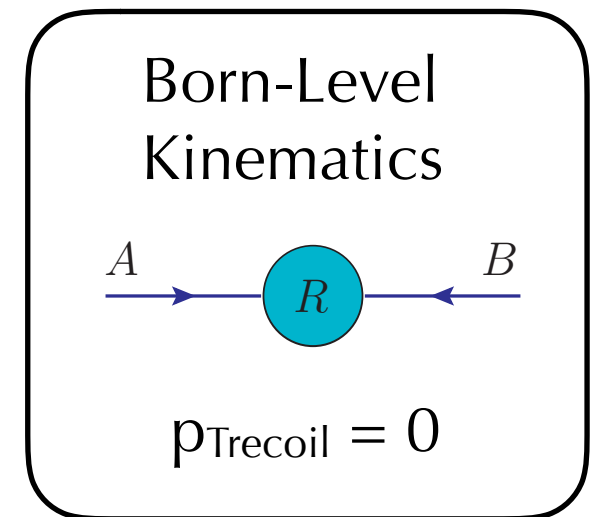
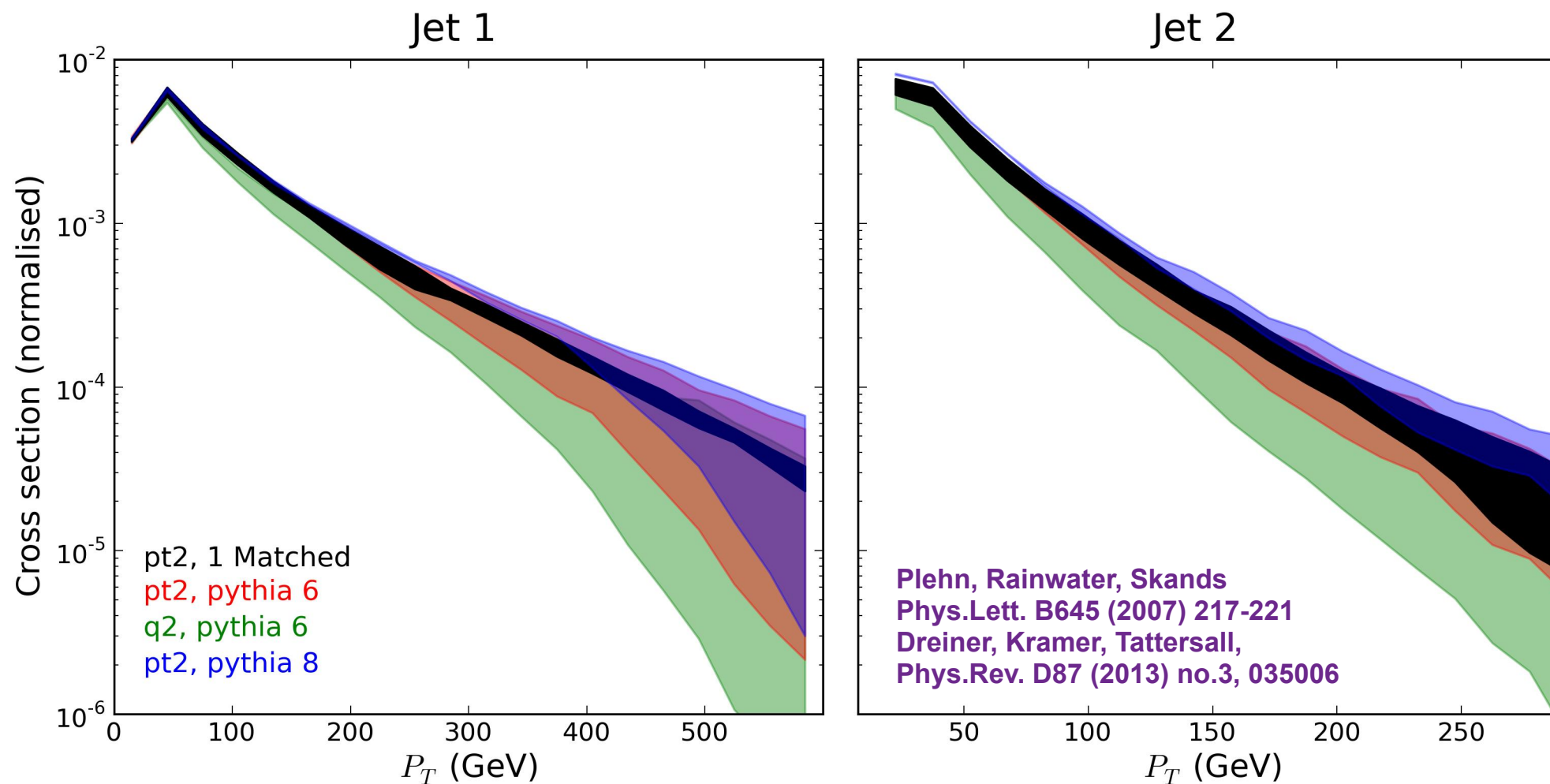
Signal Properties

(adapted from slides by S. Prestel, MC4BSM workshop)

Do you need ME + PS for BSM signals?

Example: pair production of 500-GeV squarks, plus QCD jets

Variation of shower profiles vs matrix-element-matched calc



Improved QCD pins down the transverse momenta

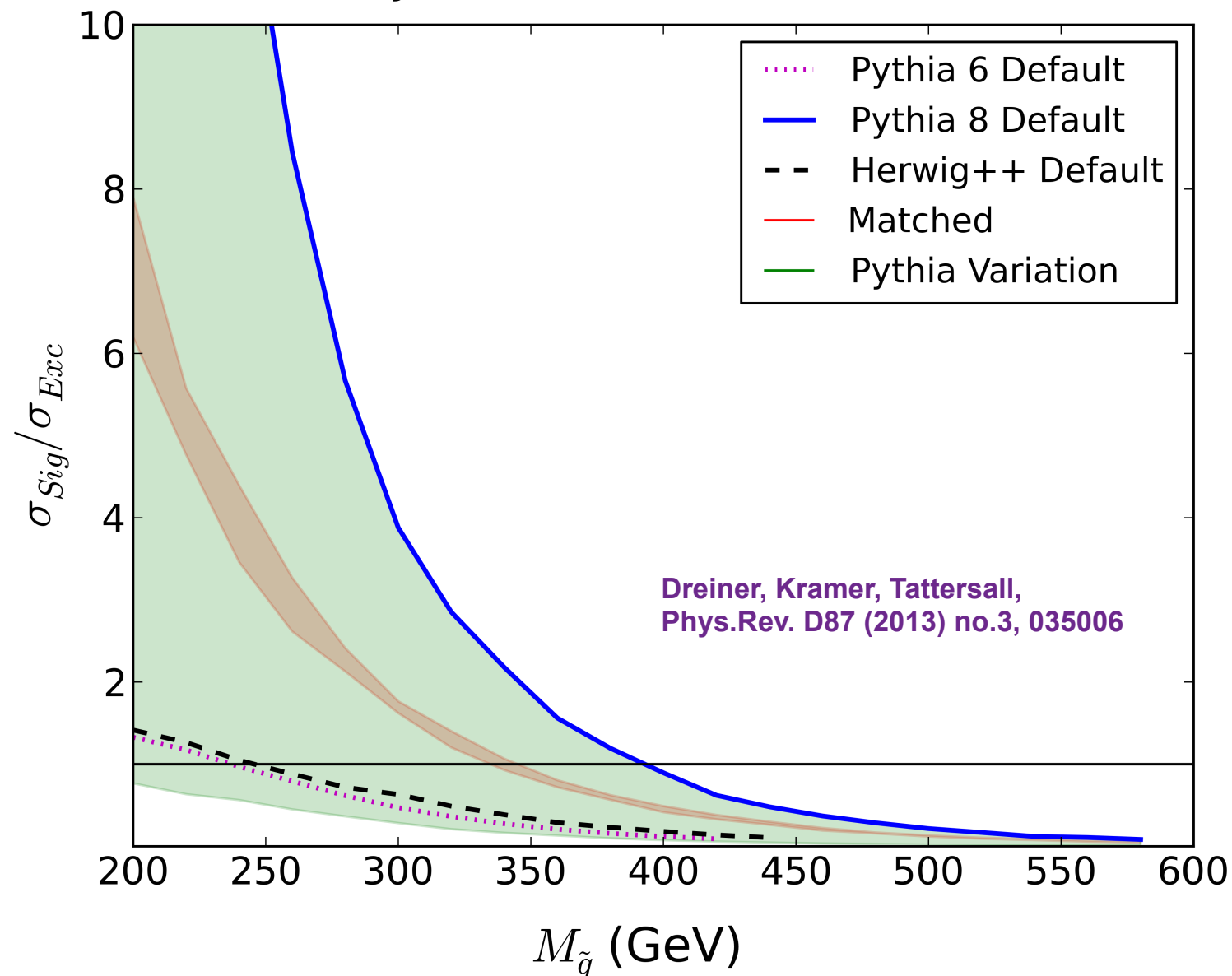
Exclusions (& Recasting)

(adapted from slides by S. Prestel, MC4BSM workshop)

How good is your exclusion?

Variation of shower profiles vs matrix-element-matched calc

Monojet Search Limits, $\sqrt{s} = 7$ TeV



+ “Recasting” is getting to be a big activity; reinterpreting high-energy collider analyses.

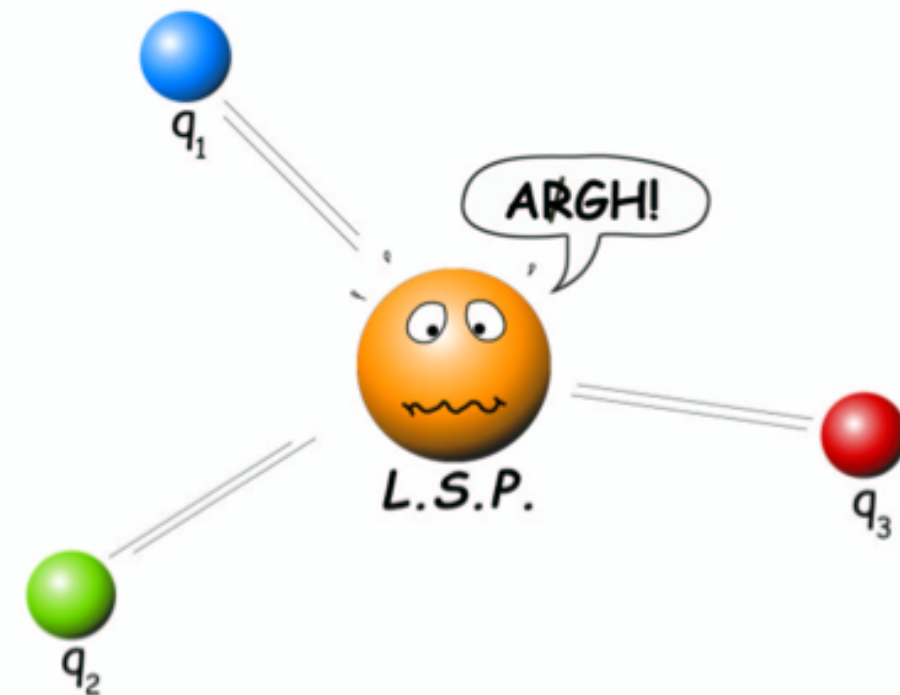
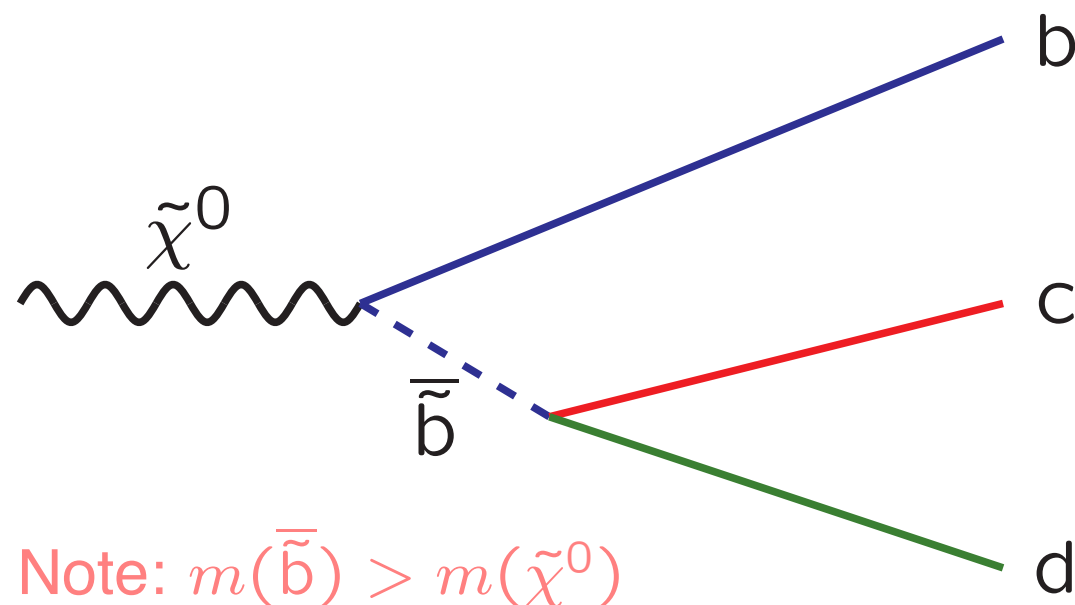
Cranmer, Yavin
JHEP 1104 (2011) 038

Basically use MC to reinterpret exclusion limits done for one model/topology in context of another model!

Exotic Colours

(adapted from slide by T. Sjostrand)

Baryon number violation (BNV) is allowed in SUSY superpotential. Alternatively lepton number violation, but proton unstable if both. BNV couplings should not be too big, or else large loop corrections \Rightarrow relevant for LSP (Lightest Supersymmetric Particle).



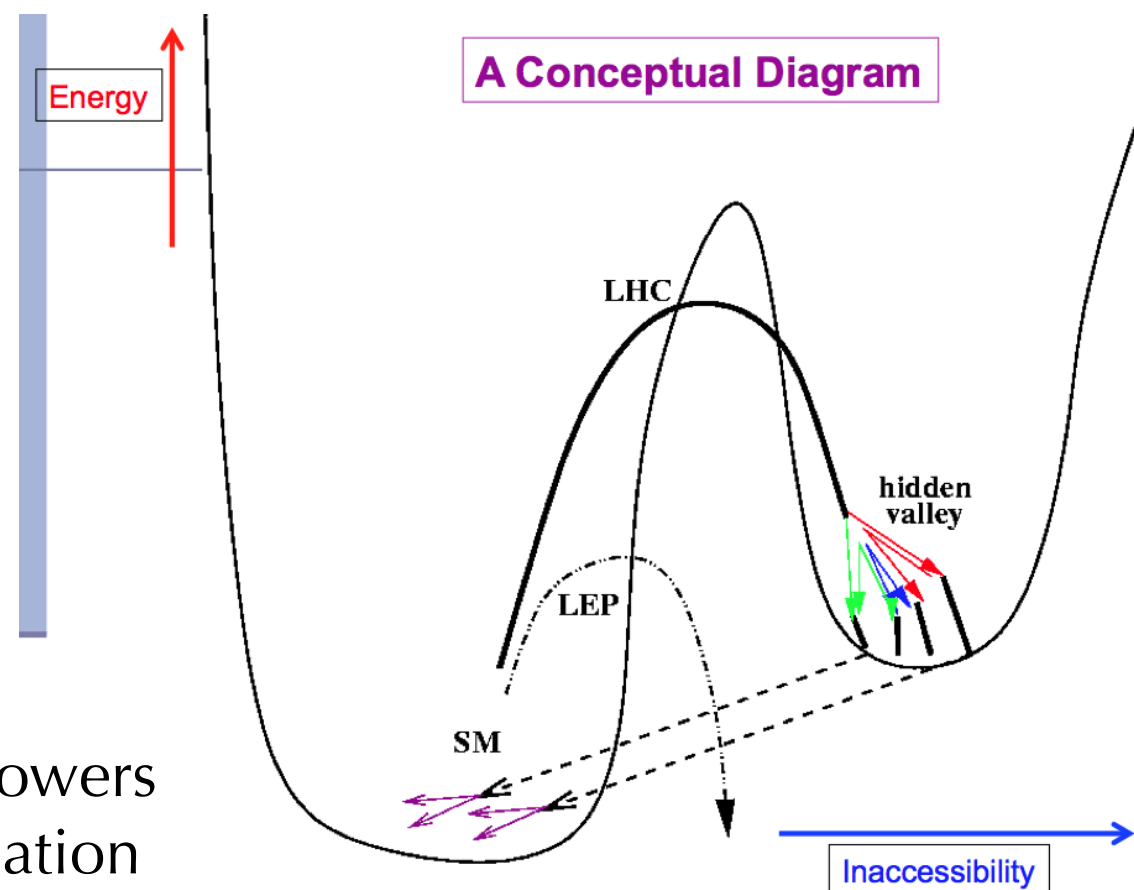
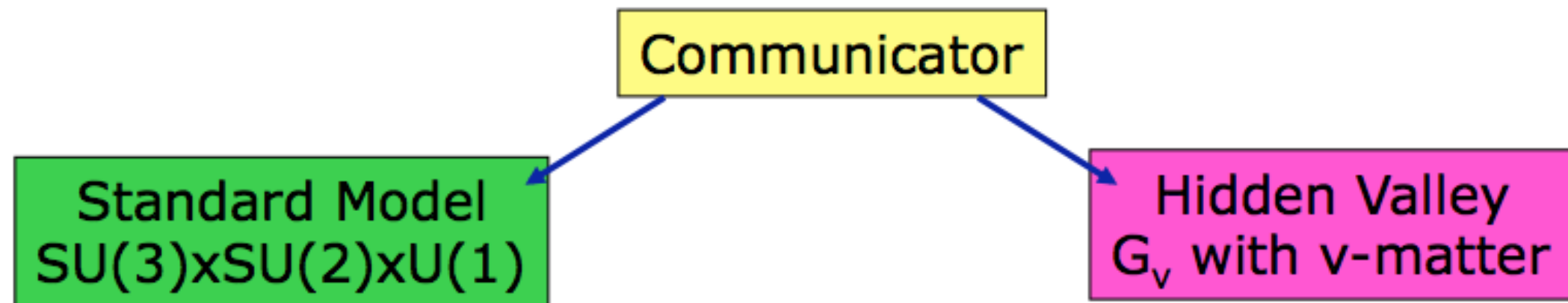
What about showers and hadronization in decays?

P. Skands & TS, Nucl. Phys. B659 (2003) 243;

N. Desai & P. Skands, arXiv:1109.5852 [hep-ph]

Hidden Valleys / Emerging Jets

M. Strassler, K. Zurek, Phys. Lett. B651 (2007) 374; ...



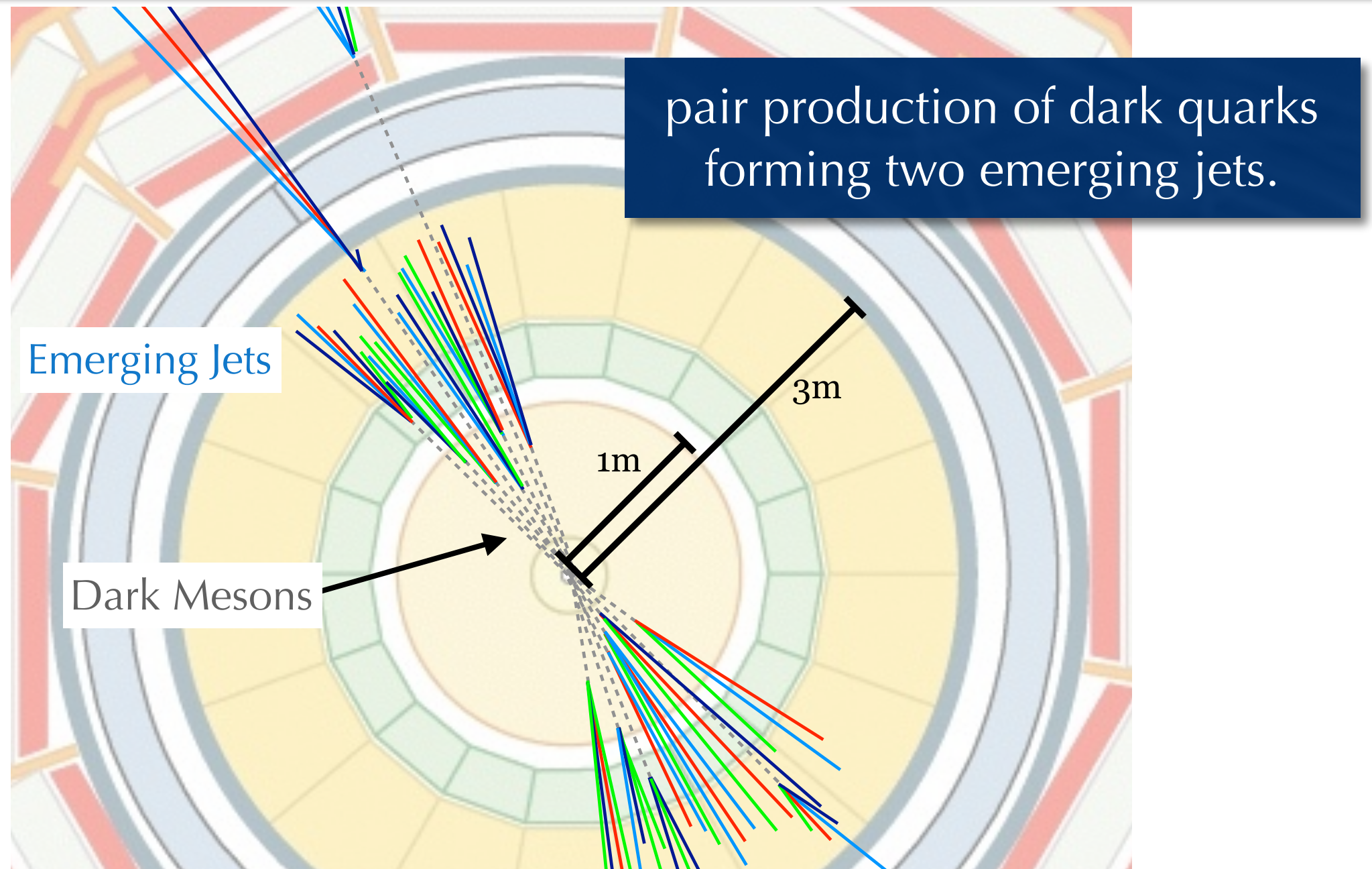
Hidden Valley
aka "Dark" Sector
aka "Hidden" Sector

Courtesy
M. Strassler

Hidden-Valley Showers
+ Valley Hadronisation

↳ L. Carloni & TS, JHEP 1009, 105; L. Carloni, J. Rathsman & TS, JHEP 1104, 091

Hidden Valleys / Emerging Jets



Requirements for a model to produce emerging jet phenomenology:

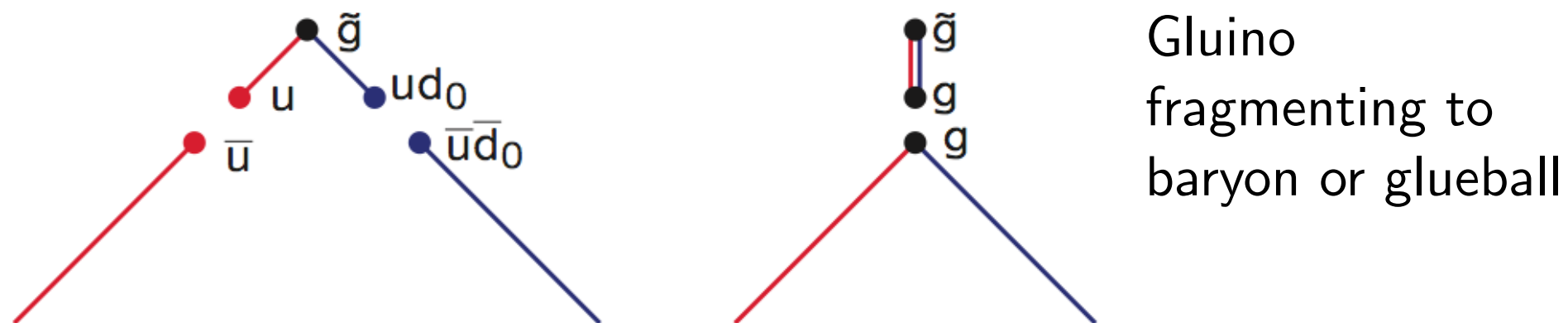
- Hierarchy between the mediator mass and hidden sector mass.
- Strong coupling in hidden sector \rightarrow large particle multiplicity.
- Macroscopic decay lengths of hidden sector fields back to the visible sector

Schwaller, Stolarski, Weiler
JHEP 1505 (2015) 059

R-Hadrons

⇒ PYTHIA allows for hadronization of 3 generic states:

- colour octet uncharged, like \tilde{g} , giving $\tilde{g}u\bar{d}$, $\tilde{g}uud$, $\tilde{g}g$, ... ,
- colour triplet charge $+2/3$, like \tilde{t} , giving $\tilde{t}\bar{u}$, $\tilde{t}ud_0$, ... ,
- colour triplet charge $-1/3$, like \tilde{b} , giving $\tilde{b}\bar{c}$, $\tilde{b}su_1$,



Most hadronization properties by analogy with normal string fragmentation, but

glueball formation new aspect, assumed $\sim 10\%$ of time (or less).

R-hadron interactions with matter: part of detector simulation, i.e. GEANT, not PYTHIA
Freight-train BSM particle surrounded by light pion/gluon cloud \rightarrow little dE/dx
+ charge flipping ! A.C. Kraan, Eur. Phys. J. C37 (2004) 91; M. Fairbairn et al., Phys. Rep. 438 (2007) 1

Interfaces

Monte Carlo for BSM often involves **chains** of codes
→ **interfaces** play a central role

Model Building Tool (e.g., FeynRules, LanHEP, ...)

→ Model File (e.g., UFO) [arXiv:1108.2040](https://arxiv.org/abs/1108.2040)

Matrix-Element Generator (e.g., MadGraph, CalcHEP, ...)

→ Matrix Elements

→ Les Houches Event Files (LHEF) [hep-ph/0609017](https://arxiv.org/abs/hep-ph/0609017)

SLHA for SUSY spectra: [hep-ph/0311123](https://arxiv.org/abs/hep-ph/0311123) [arXiv:0801.0045](https://arxiv.org/abs/hep-ph/0801.0045)

BSM-SLHA for new particles: [arXiv:0712.3311](https://arxiv.org/abs/hep-ph/0712.3311)

Matching/Merging Strategy + Parton Shower Generator

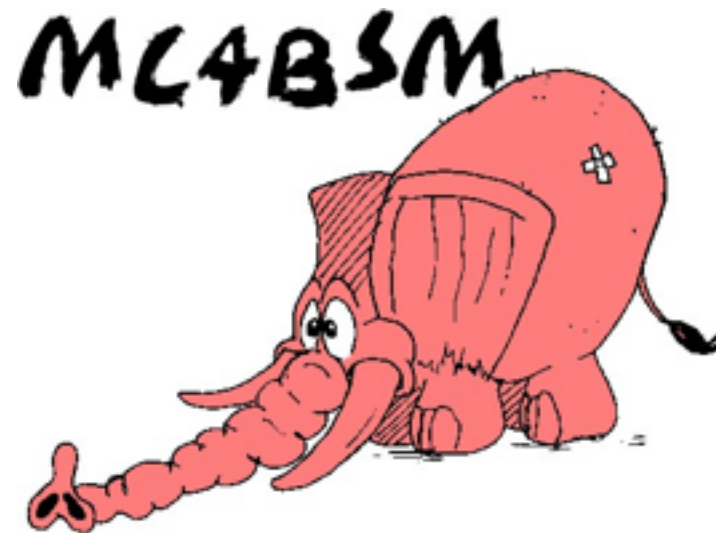
→ Interpretation of LHEF

→ Hadron-Level event files (HepMC) → Analysis Tools
(e.g., FASTJET, DELPHES, ...)

MC for BSM

Yearly workshop series **MC4BSM** (running since 2006)

“Gather theorists and experimentalists interested in developing Monte Carlo tools to simulate signatures of BSM Physics, and to use such tools in searches and phenomenological studies”




<http://theory.fnal.gov/mc4bsm/>

This year, 20-24 July, UCAS-YuQuan (Beijing, China)

<http://indico.ihep.ac.cn/event/5301/>

Next year (2017): **SLAC** (Stanford, California)

The End

A complex visualization of particle tracks from a proton-proton collision. The tracks are represented by a dense web of thin green lines radiating from a central point. A few thicker, colored lines (blue, yellow, red) stand out among the green ones. The background is black, making the green lines highly visible.

THANK YOU

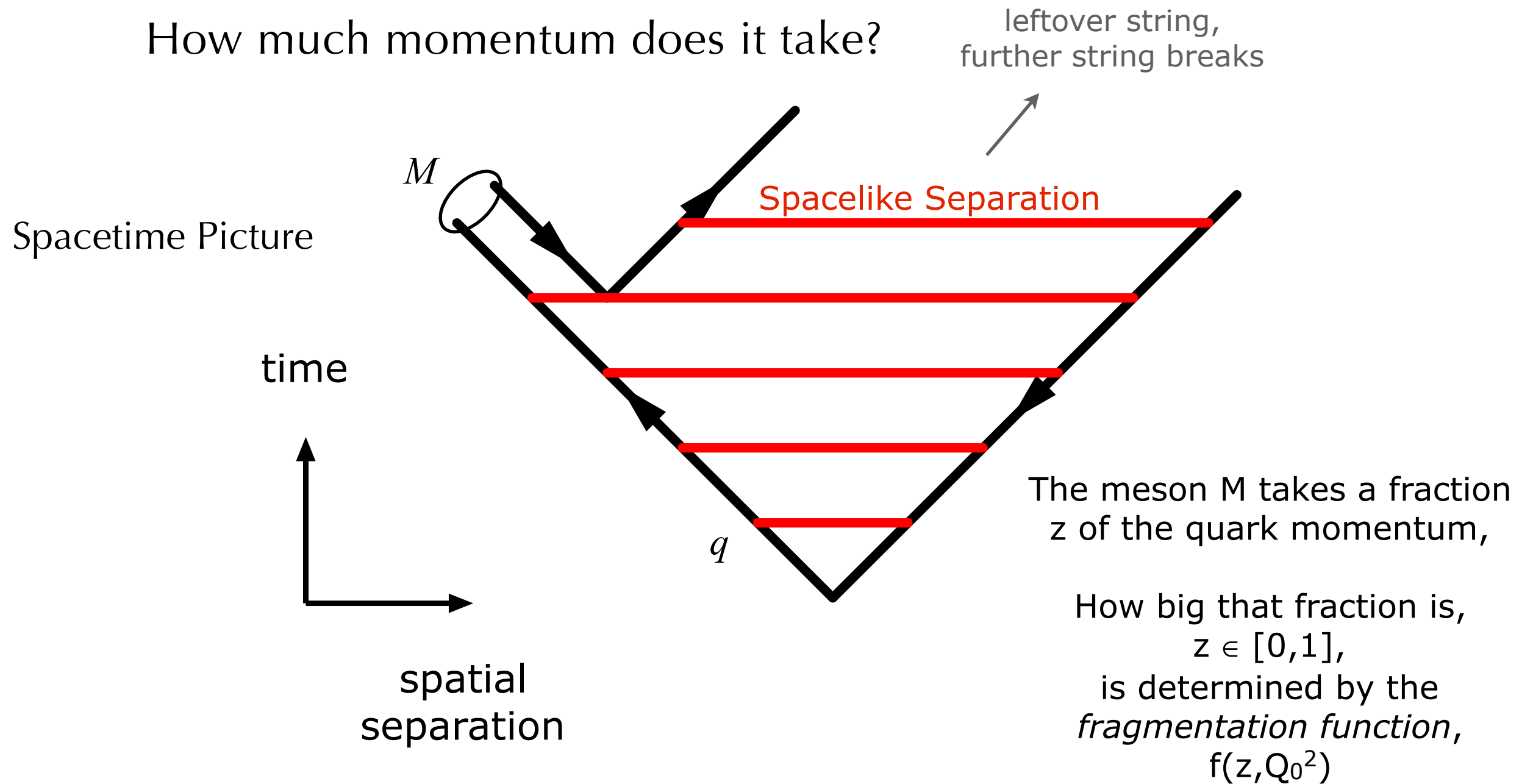
(Event recorded by CMS with 78 proton-proton collisions)

Extra Slides

Fragmentation Function

Having selected a hadron flavor

How much momentum does it take?

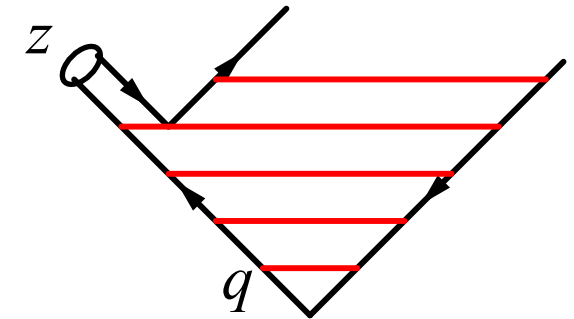


Left-Right Symmetry

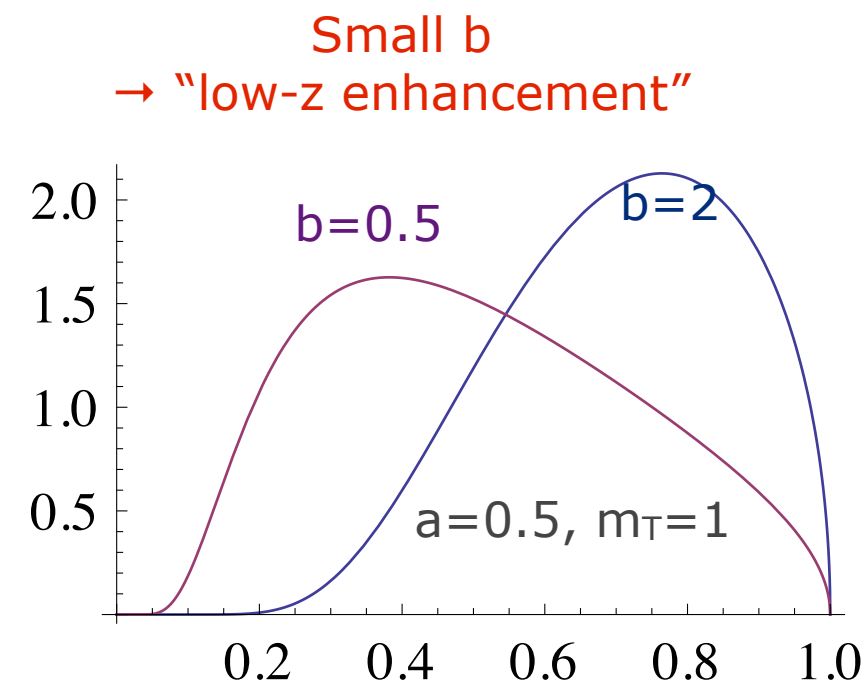
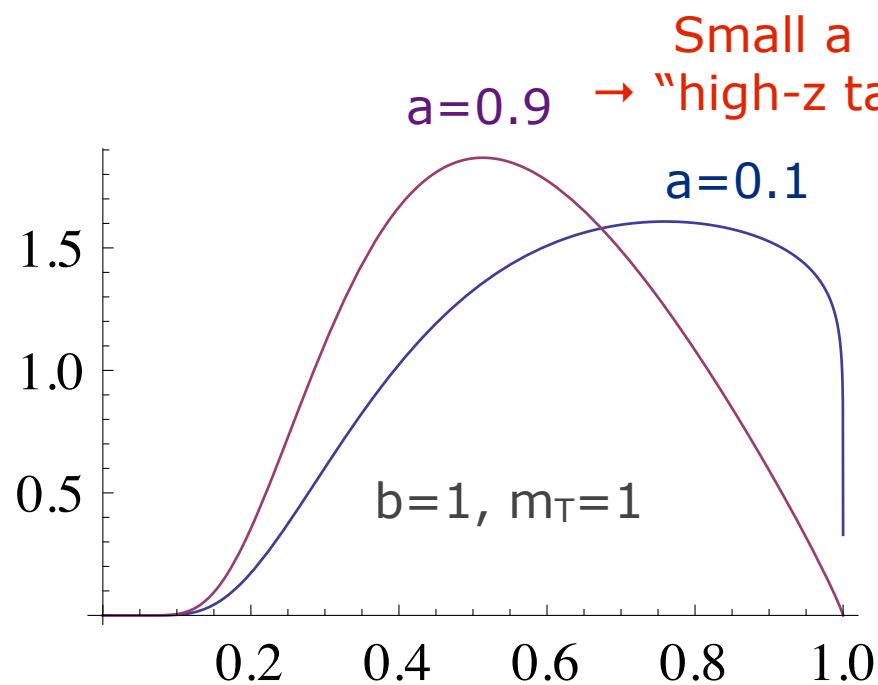
Causality → Left-Right Symmetry

→ Constrains form of fragmentation function!

→ Lund Symmetric Fragmentation Function



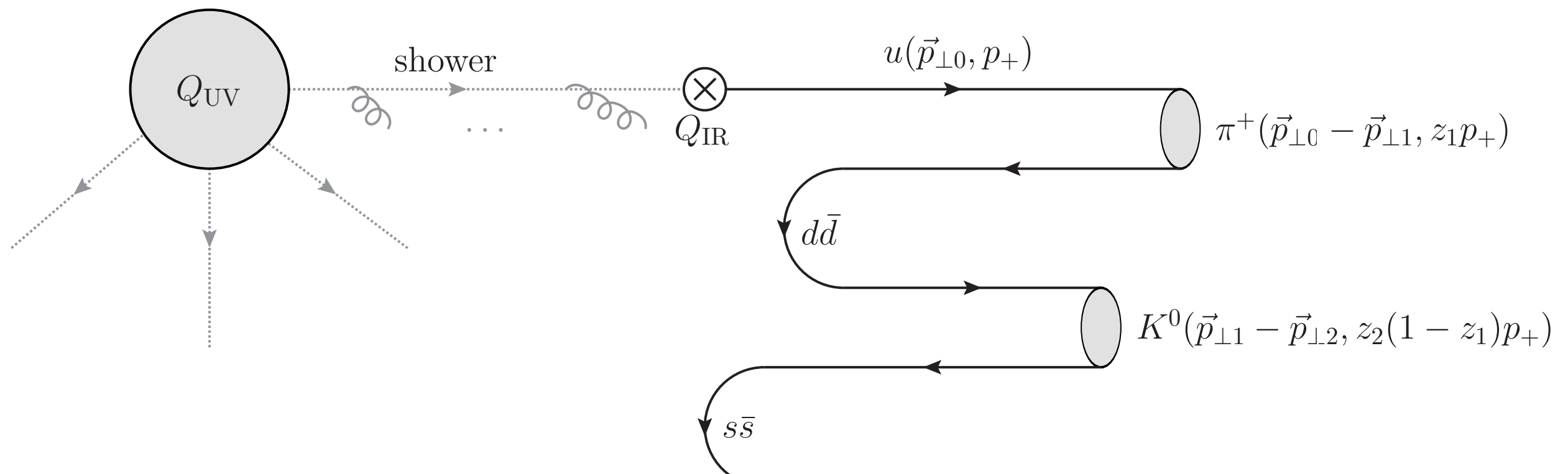
$$f(z) \propto \frac{1}{z} (1-z)^a \exp\left(-\frac{b(m_h^2 + p_{\perp h}^2)}{z}\right)$$



Note: In principle, a can be flavour-dependent. In practice, we only distinguish between baryons and mesons

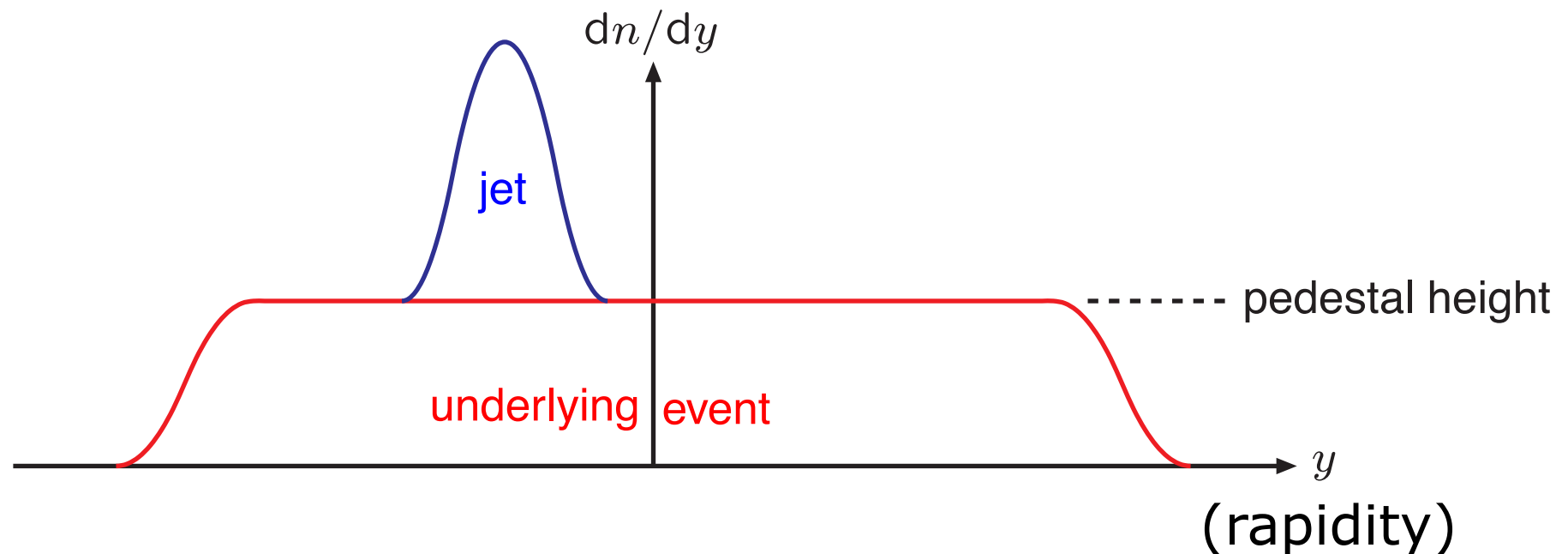
Iterative String Breaks

Causality → May iterate from outside-in



What is Underlying Event ?

“Pedestal Effect”



Useful variable in hadron collisions: **Rapidity**

Designed to be additive
under Lorentz Boosts along
beam (z) direction

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$$

$$y \rightarrow -\infty \text{ for } p_z \rightarrow -E$$

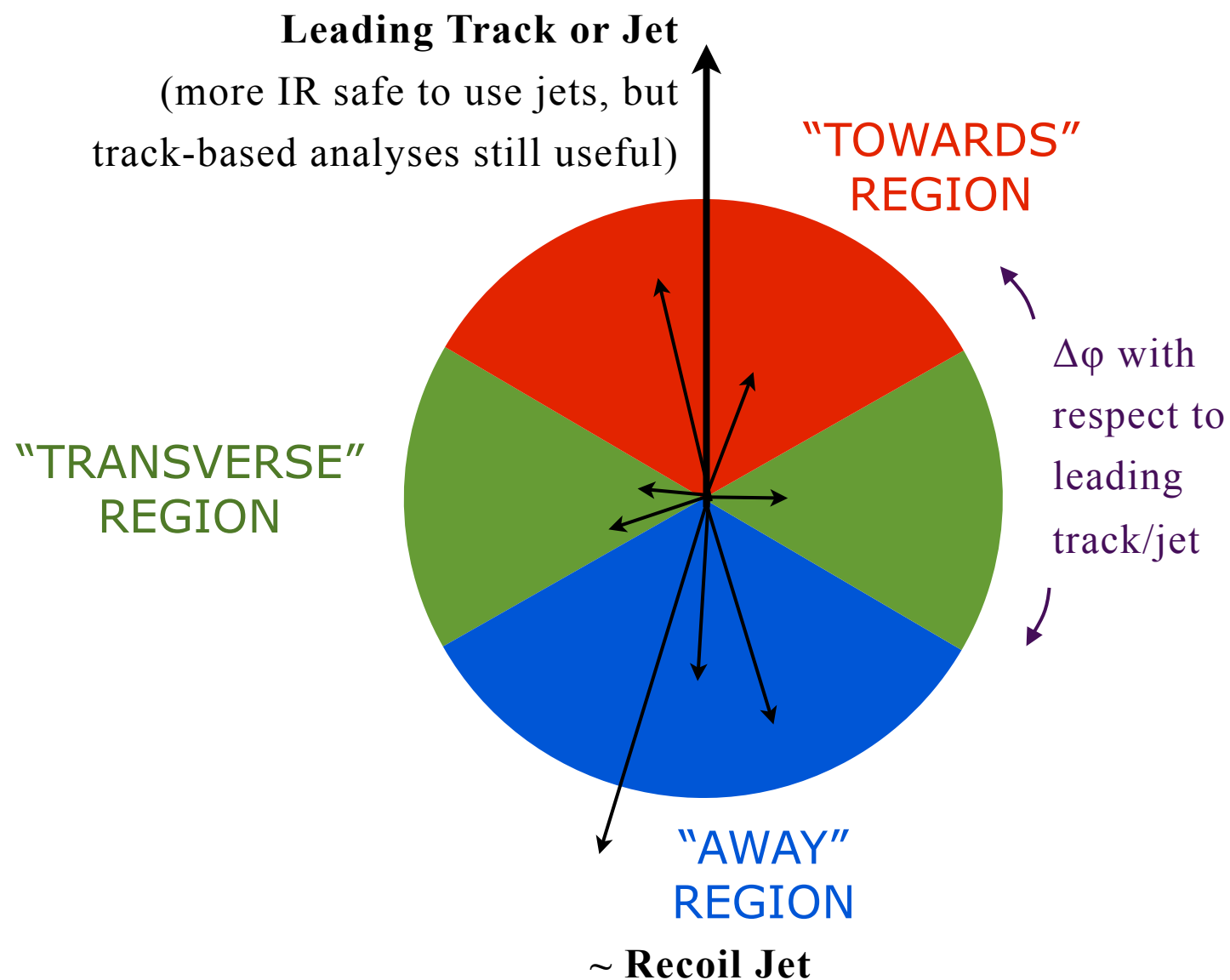
$$y \rightarrow 0 \text{ for } p_z \rightarrow 0$$

$$y \rightarrow \infty \text{ for } p_z \rightarrow E$$

Illustrations by T. Sjöstrand

The "Rick Field" UE Plots

There are many UE variables.
The most important is $\langle \Sigma p_T \rangle$ in the "Transverse Region"

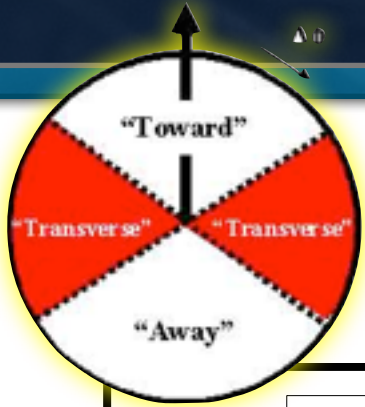


Transverse Region (TRNS)

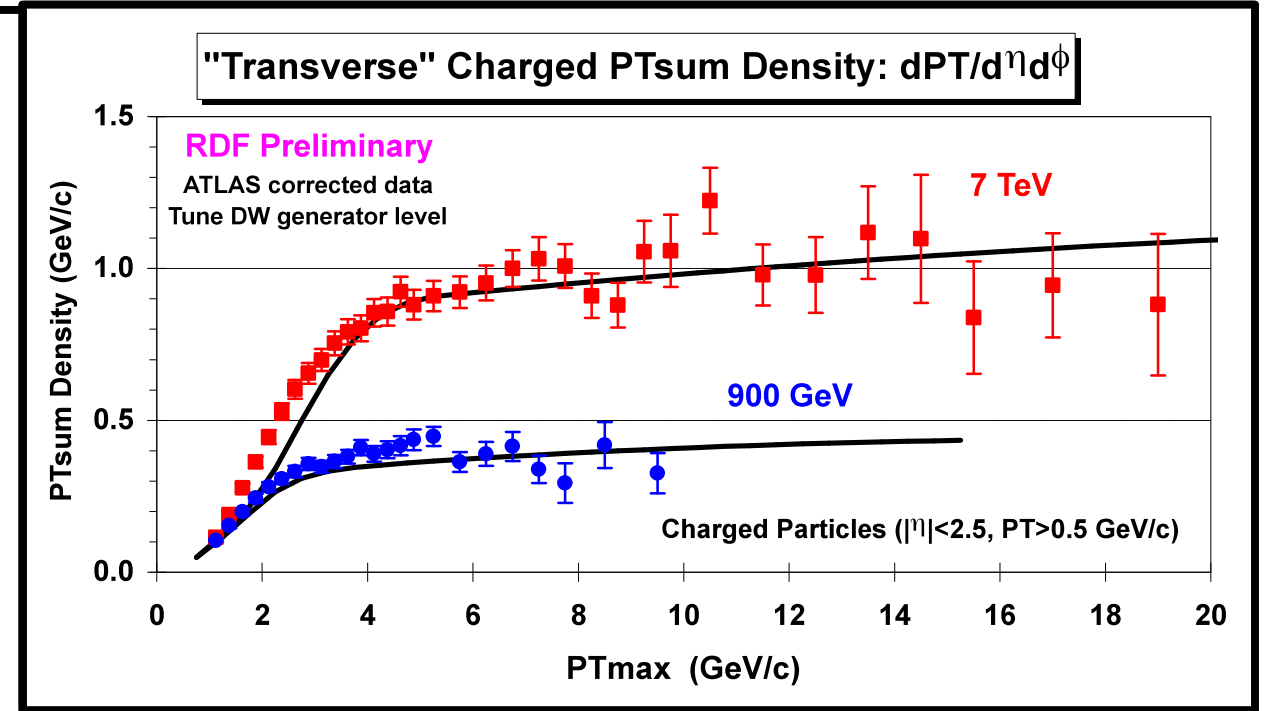
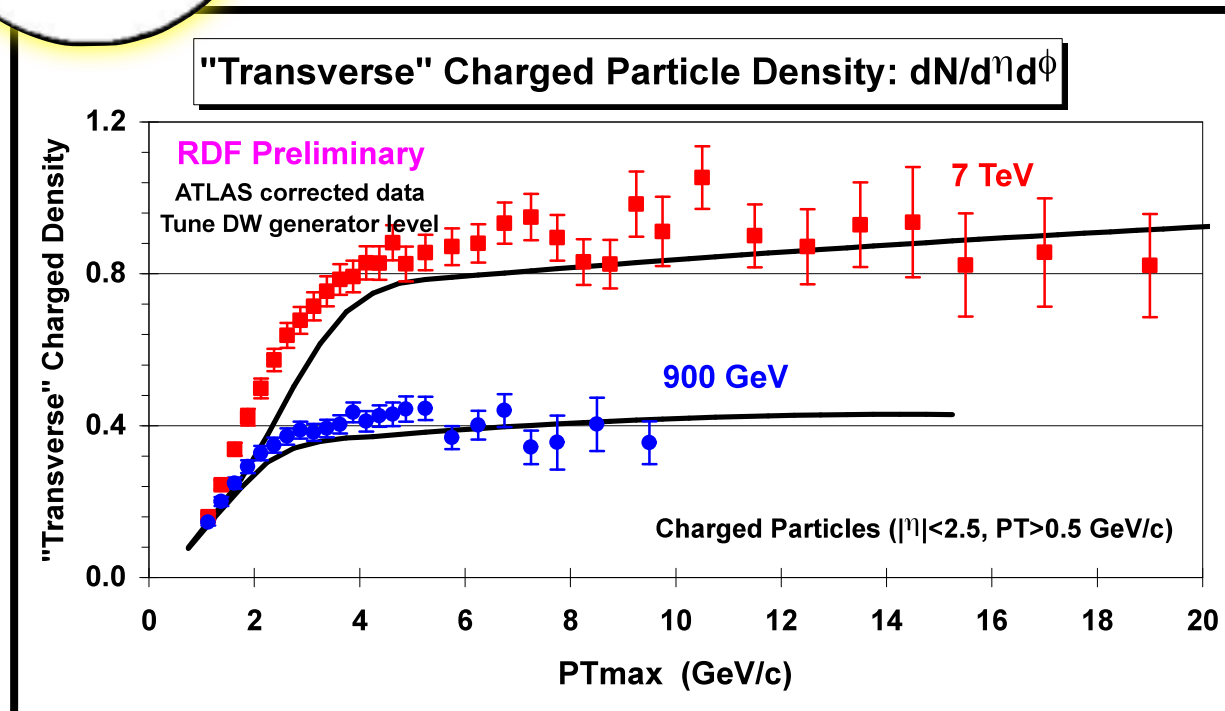
Sensitive to activity at right angles to the hardest jets

Useful definition of Underlying Event

The Pedestal



LHC from 900 to 7000 GeV - ATLAS



Track Density (TRANS)

Not Infrared Safe

Large Non-factorizable Corrections

Prediction off by $\approx 10\%$

Sum(pT) Density (TRANS)

(more) Infrared Safe

Large Non-factorizable Corrections

Prediction off by $< 10\%$

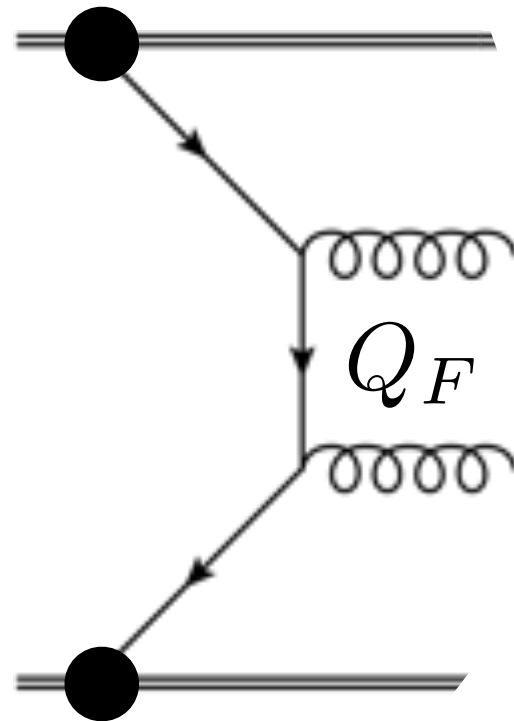
Truth is in the eye of the beholder:

R. Field: "See, I told you!"

Y. Gehrstein: "they have to fudge it again"

Physics of the Pedestal

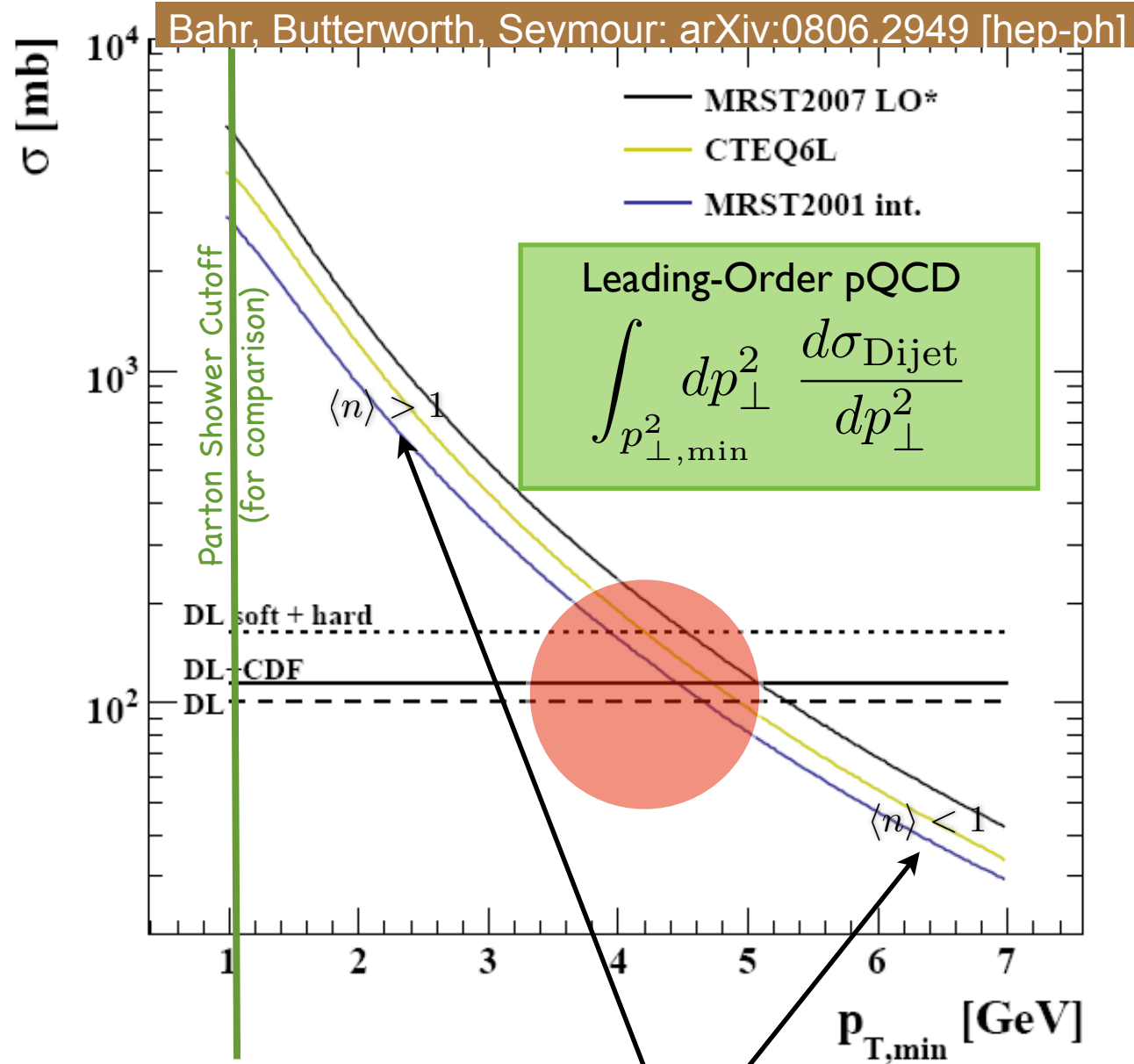
Factorization: Subdivide Calculation



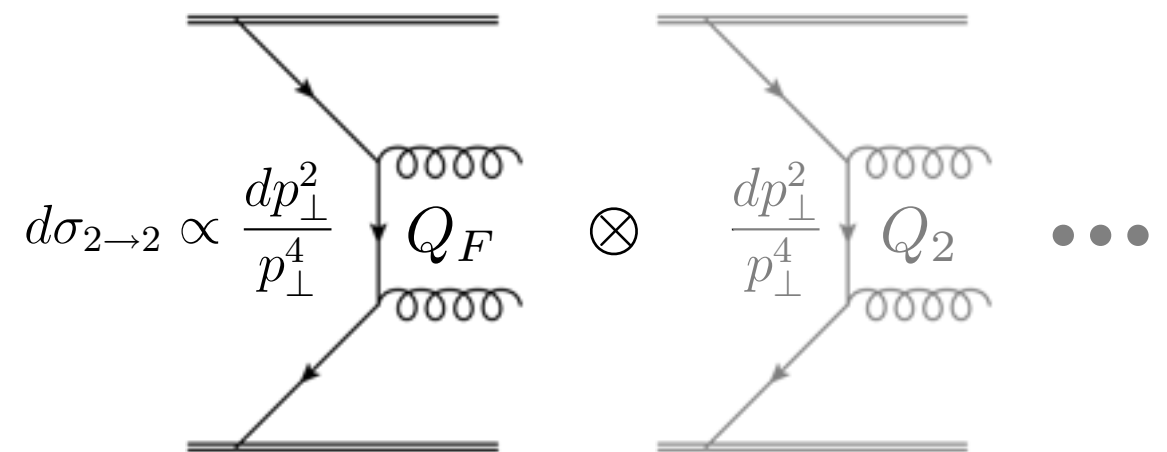
- Multiple Parton Interactions* go beyond existing theorems
- perturbative short-distance physics in Underlying Event
 - Need to generalize factorization to MPI

Multiple Parton Interactions

= Allow several parton-parton interactions per hadron-hadron collision. Requires extended factorization ansatz.



Earliest MC model ("old" PYTHIA 6 model)
Sjöstrand, van Zijl PRD36 (1987) 2019



Lesson from bremsstrahlung in pQCD:
divergences \rightarrow fixed-order breaks down
Perturbation theory still ok, with
resummation (unitarity)

\rightarrow Resum dijets?
Yes \rightarrow MPI!

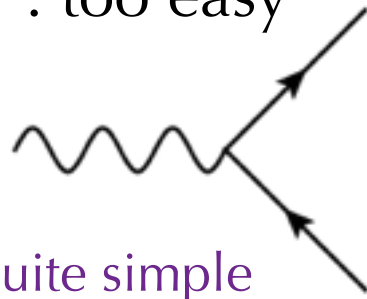
$$\sigma_{2 \rightarrow 2}(p_{\perp \min}) = \langle n \rangle(p_{\perp \min}) \sigma_{\text{tot}}$$

Parton-Parton Cross Section Hadron-Hadron Cross Section

Colour Confusion

Between which partons do confining potentials arise?

e^+e^- : too easy



(still quite simple even after including bremsstrahlung etc.)

At e^+e^- colliders (eg LEP) : generally good agreement between **measured** particle spectra and **models** based on parton/antenna showers + strings

Basically a single **3-3bar** system, very close to the original lattice studies motivating the string model.

(+ extensions to WW reasonable to $\sim O(1/N_c^2)$)

→ re-use same models as input for LHC (universality) ?

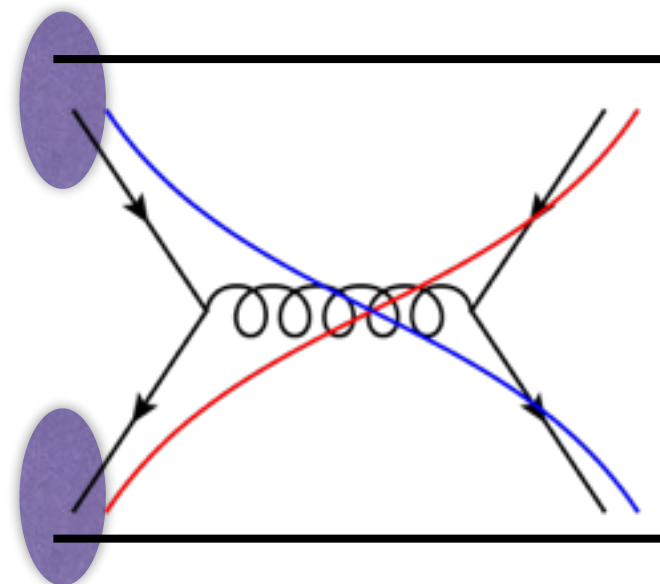
Proton-Proton (LHC)

A lot more colour kicked around (& also colour in initial state)

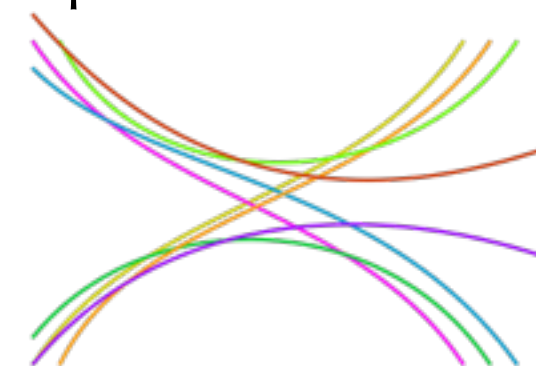
Include "Beam Remnants"

Still might look relatively simple, to begin with

(+baryon beam remnants → "string junctions")



But no law against *several* parton-parton interactions



In fact, can easily be shown to happen frequently
Included in all (modern) Monte Carlo models
But how to make sense of the colour structure?

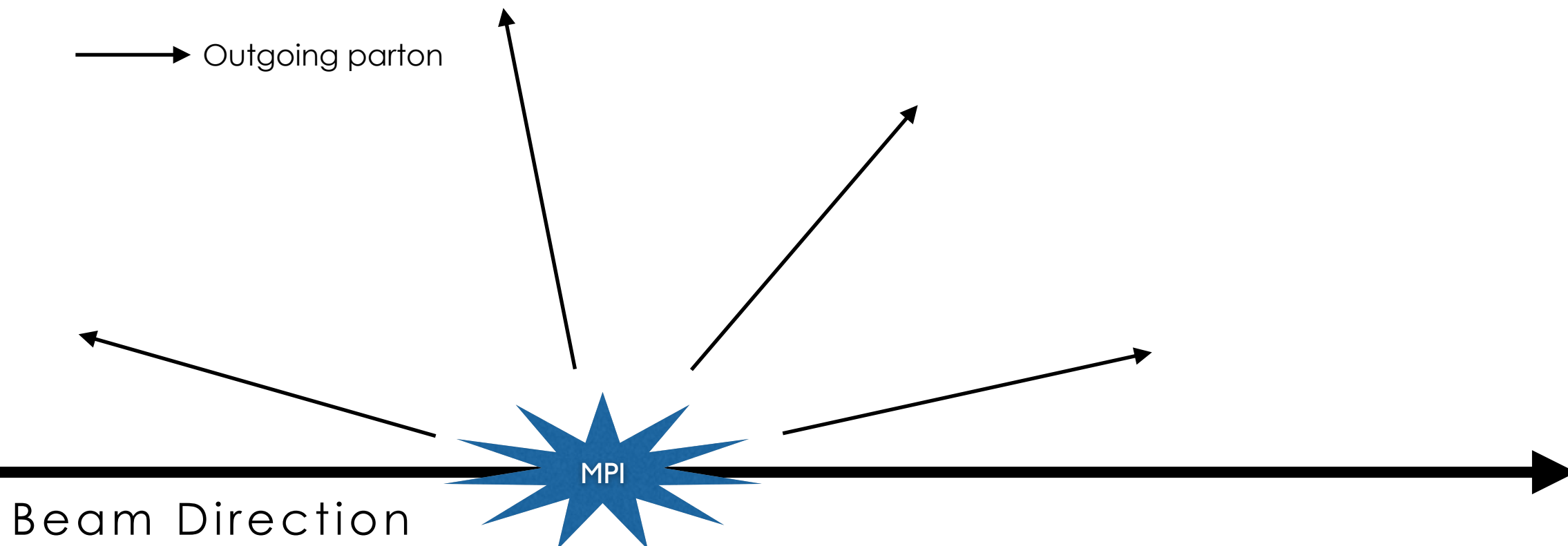
String-fragmentation of junctions: Sjöstrand & Skands **Nucl.Phys. B659 (2003) 243**

Colour: What's the Problem?

(including **MPI**: Multiple Parton-Parton Interactions ~ the “underlying event”)

Without Colour Reconnections

Each MPI hadronizes **independently** of all others



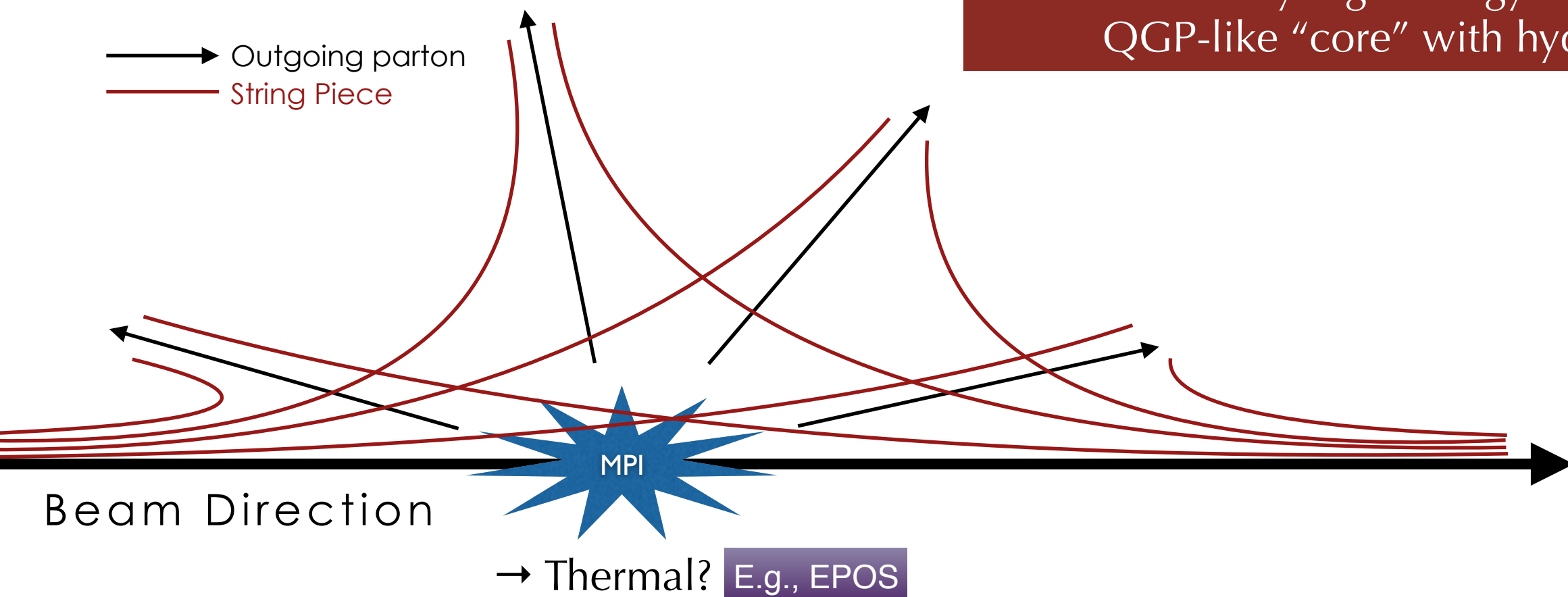
Colour: What's the Problem?

(including **MPI**: Multiple Parton-Parton Interactions ~ the “underlying event”)

Without Colour Reconnections

Each MPI hadronizes **independently** of all others

So many strings in so little space
If true → Very high energy densities
QGP-like “core” with hydro?



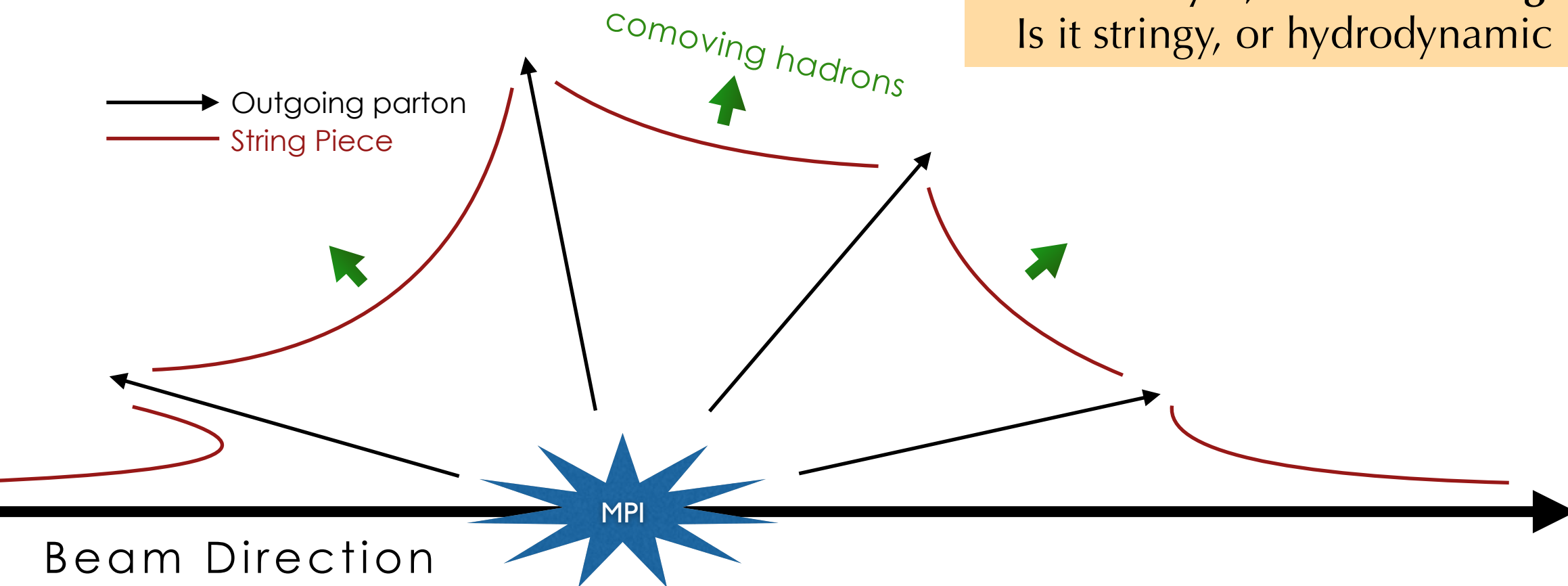
Colour Reconnections

(including **MPI**: Multiple Parton-Parton Interactions ~ the “underlying event”)

With Colour Reconnections
MPI hadronize **collectively**

See also Ortiz et al., Phys.Rev.Lett. 111 (2013) 4, 042001

Highly interesting theory questions now.
Is there collective flow in pp? Or not?
If yes, what is its origin?
Is it stringy, or hydrodynamic ? (or ...?)



String-Length Minimisation E.g., PYTHIA, HERWIG

Or Thermal? E.g., EPOS

Or Higher String Tension? E.g., DIPSY rope

What are “Colour Reconnections”?

Simple example: $e^+e^- \rightarrow W^+W^- \rightarrow \text{hadrons}$

Intensely studied at LEP2.

CR implied a non-perturbative uncertainty on the W mass measurement, $\Delta M_W \sim 40 \text{ MeV}$

CR constrained to $\sim 10\% \sim 1/N_C^2$

Simple two-string system. What about pp?

Several modelling attempts

Based on “just” minimising the string action

String interactions (Khoze, Sjostrand)

Generalized Area Law (Rathsman et al.)

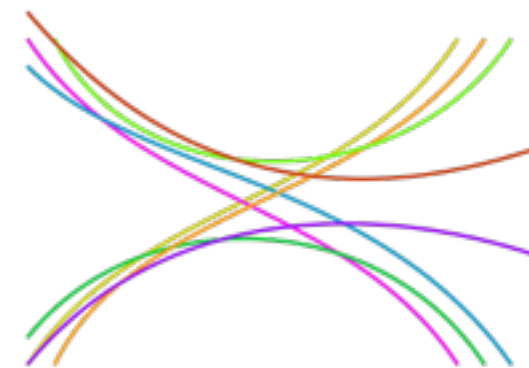
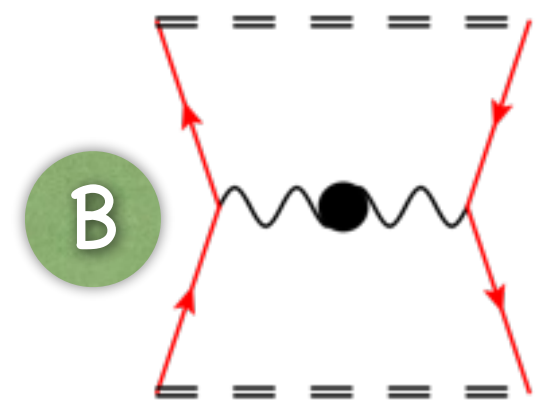
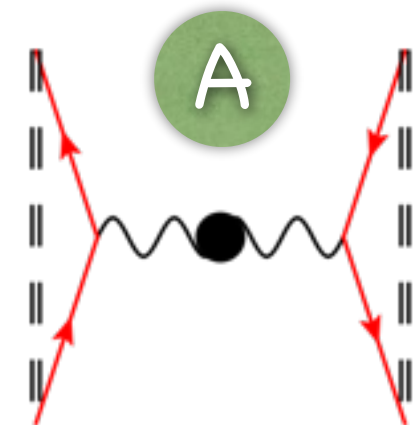
Colour Annealing (Skands et al.)

Gluon Move Model (Sjostrand et al.)

More recently: $SU(3)_C$ group multiplet weights

Dipole Swing (Lonnblad et al.)

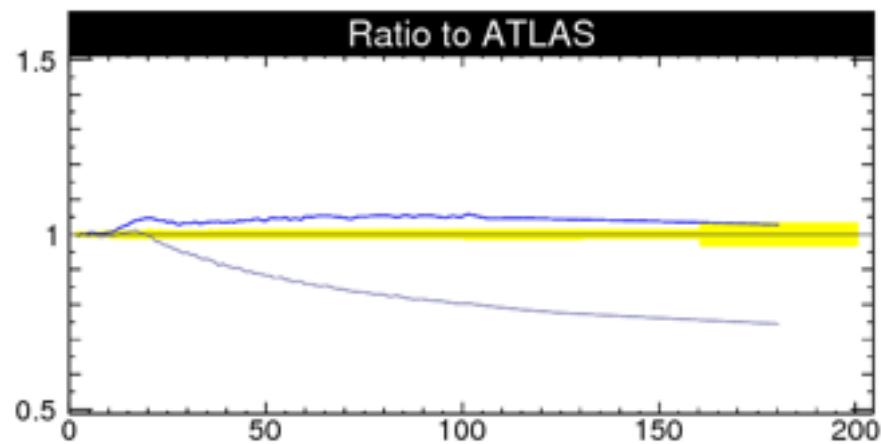
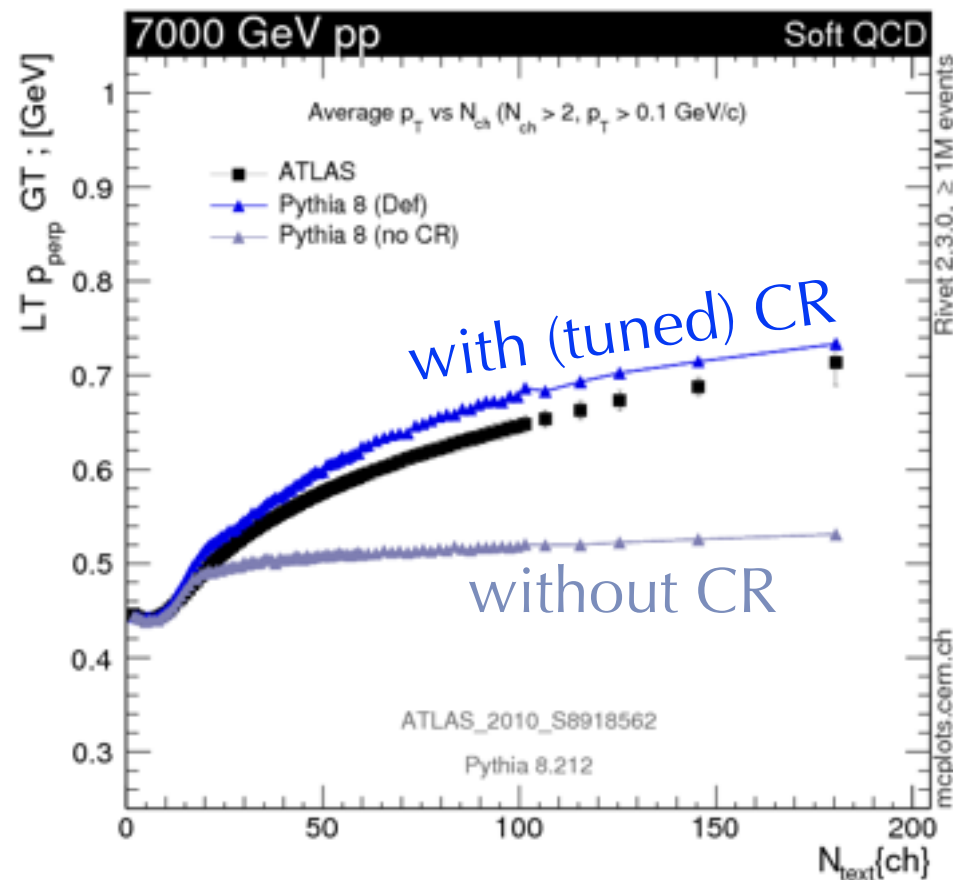
String Formation Beyond Leading Colour (Skands et al.)



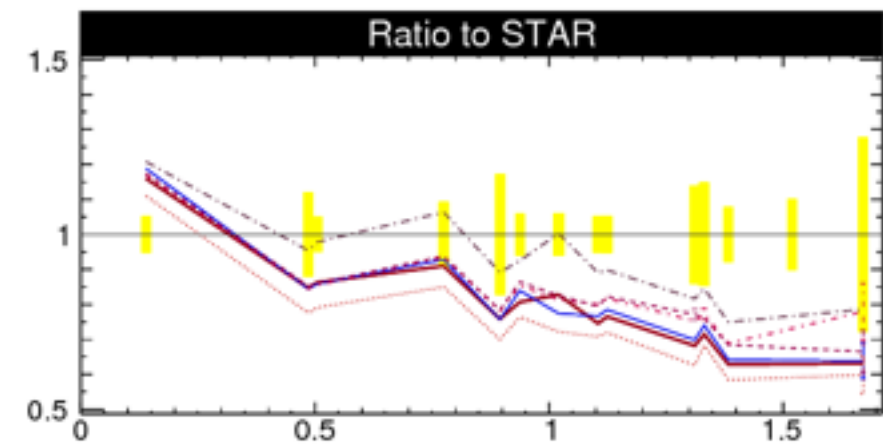
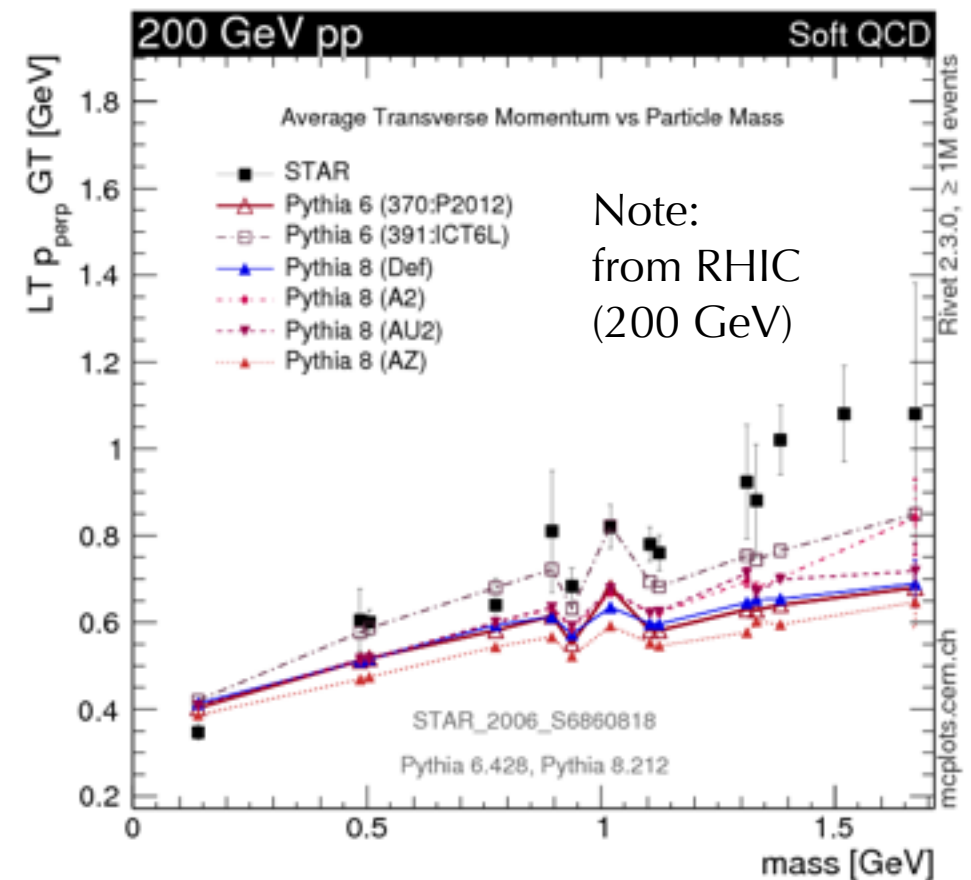
$$\begin{aligned}
 3 \otimes \bar{3} &= 8 \oplus 1 \\
 3 \otimes 3 &= 6 \oplus \bar{3} \\
 3 \otimes 8 &= 15 \oplus 6 \oplus 3 \\
 8 \otimes 8 &= 27 \oplus 10 \oplus \bar{10} \oplus 8 \oplus 8 \oplus 1
 \end{aligned}$$

What do we see?

$\langle p_T \rangle$ vs Number of Particles



$\langle p_T \rangle$ vs Particle Mass



Average p_T increases with particle multiplicity and (faster than predicted) with particle mass